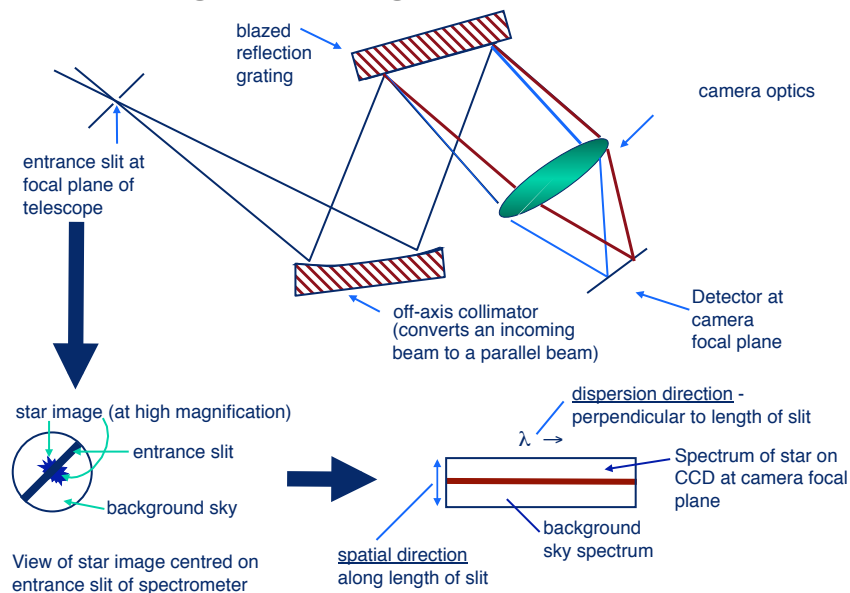


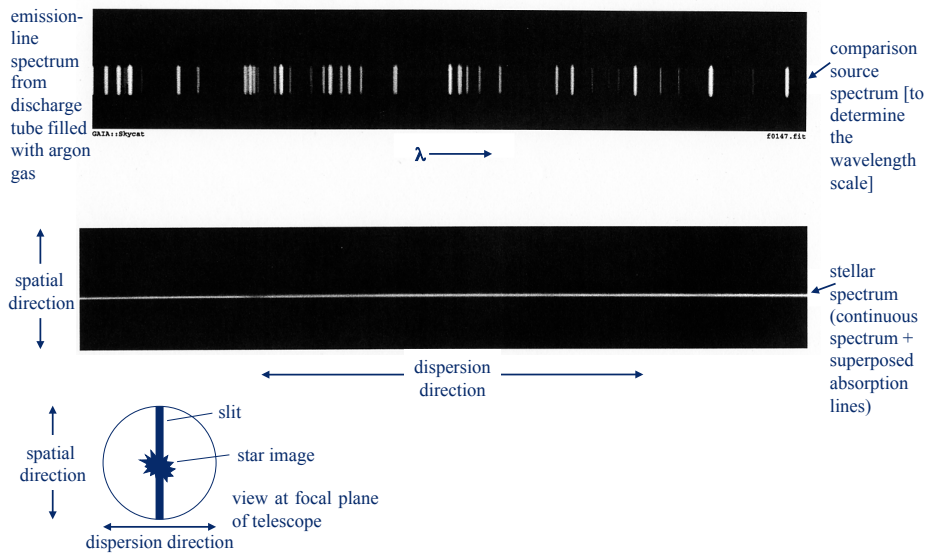
10. Spectroscopy

- **Slit spectroscopy: entrance aperture of the spectrometer is a narrow slit.**
- **Entrance slit must be**
 - wide enough to allow light from the source through to spectrograph,
 - narrow enough to preserve a desired spectral resolution.
- **Typical slit width ~ 1 arcsec, while seeing discs are usually >1 arcsec, so lose some light from source.**
- **Need to design spectrograph optics to give required spectral resolution for a given slit width.**
 - Cannot use a slit spectrograph for spectrophotometry because light losses vary with seeing, unless you open the slit wide.

Grating spectrograph:schematic

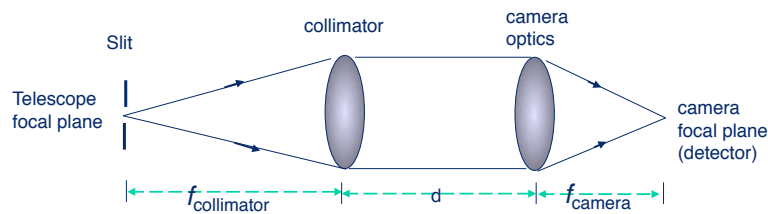


'Raw' spectra recorded by a CCD on a spectrometer



Collimator, camera and image size

- **Image scale at slit (telescope focal plane):**
 $206265/f_t \text{ arcsec mm}^{-1}$
- **Image of slit is projected onto CCD in spectrograph camera focal plane by 2 optical elements: collimator and camera.**
 - Need to match projected slit width to CCD pixel size.



- Think of *collimator* as telescope of focal length $f_{\text{collimator}}$ acting in reverse.

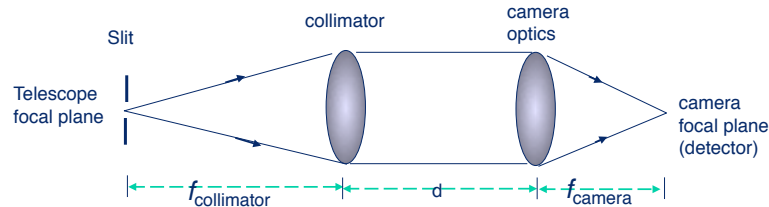
- Slit size x_{slit} is related to spread θ in beam directions reflected off collimator by: $x_{\text{slit}} = f_{\text{collimator}} \theta$.

- The *camera* is a “telescope” of focal length f_{camera}

- Size of slit image on CCD is $x_{\text{CCD}} = f_{\text{camera}} \theta$.

- Hence image of slit projected on CCD is reduced in size by the magnification factor

$$M_T = \frac{x_{\text{CCD}}}{x_{\text{slit}}} = \frac{f_{\text{camera}}}{f_{\text{collimator}}}$$



Example: Leslie Rose Spectrograph

- mirror diameter $D = 50 \text{ cm}$; Cass. focus $f/11$.

- \therefore Telescope focal length $f_{\text{telescope}} = 5500 \text{ mm}$

- \therefore image scale = $206265 / f_{\text{telescope}} = 37.5 \text{ arcsec/mm}$

- seeing disc at St Andrews is $\sim 2\text{-}2.5''$ on the sky

- (narrow!) physical width of slit = $53 - 67 \mu\text{m}$

- Spectrograph collimator mirror must also be $f/11$

- $D_{\text{collimator}} = 50 \text{ mm}$

- $\therefore f_{\text{collimator}} = 550 \text{ mm}$

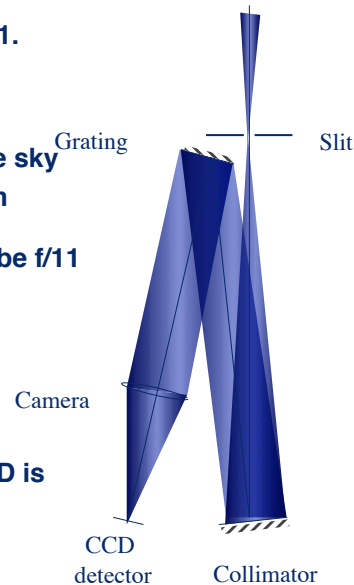
- Camera focal length $f_{\text{camera}} = 208 \text{ mm}$

- $\therefore M_T = 208/550 = 0.378$

- Hence width of slit image projected on CCD is

- $53 \text{ to } 67 \mu\text{m} \times M_T = 20 \text{ to } 25 \mu\text{m}$

- the typical sizes of CCD pixels, as required.

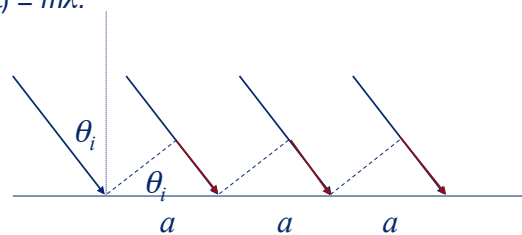


Diffraction Gratings

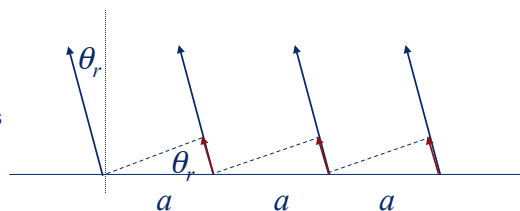
- a repeating pattern of diffracting elements, apertures or obstacles.
 - reflection diffraction gratings common because *achromatic*
 - » parallel rules scratched in an aluminium surface on a plane glass plate at hundreds of rules per mm! (& plastic copies)
 - these act as reflection phase gratings since the interference pattern is caused by changes in phase of the light waves.
- Standard grating equation: $a \sin \theta_r = m\lambda$
 - for grating spacing a (distance between adjacent rules - e.g. $a \sim 0.001$ mm) & incident beam perpendicular to the grating.
 - Zeroth order ($m = 0$) is reflected straight back ($\theta_r=0$). First order ($m = \pm 1$) gives a different θ_r , clearly dependent on λ , a .

More generally, for incident angle $\theta_i \neq 0$ the grating equation is modified to: $a (\sin \theta_i + \sin \theta_r) = m\lambda$:

Incident beam:
Path difference between adjacent grooves = $a \sin \theta_i$



Reflected beam:
Additional path difference between adjacent grooves = $a \sin \theta_r$

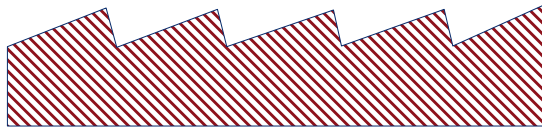


Total path difference between adjacent grooves = $a (\sin \theta_i + \sin \theta_r)$ - if diffracted beam is on same side of grating normal as incoming beam.

Otherwise path difference = $a (\sin \theta_i - \sin \theta_r)$.

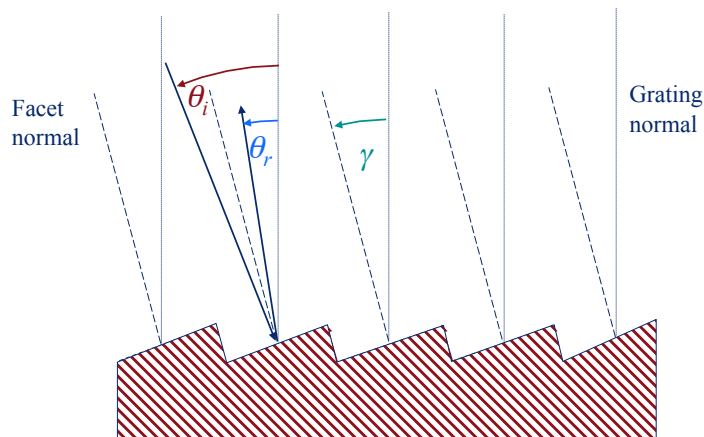
Interference

- **Outgoing beams from adjacent grooves interfere constructively when path lengths differ by an integer number of wavelengths:**
 - $m = \pm 1, \pm 2, \dots$ is constructive but $\pm 1/2, \pm 3/2 \dots$ is destructive
 - for given m , θ , depends on λ so each order is spread into a spectrum.
- **but such a grating reflects most of the light into $m = 0$ order where all wavelengths overlap - useless for spectroscopy.**
 - solution is a *blazed* reflection grating
 - grooves: inclined wedges in plane surface

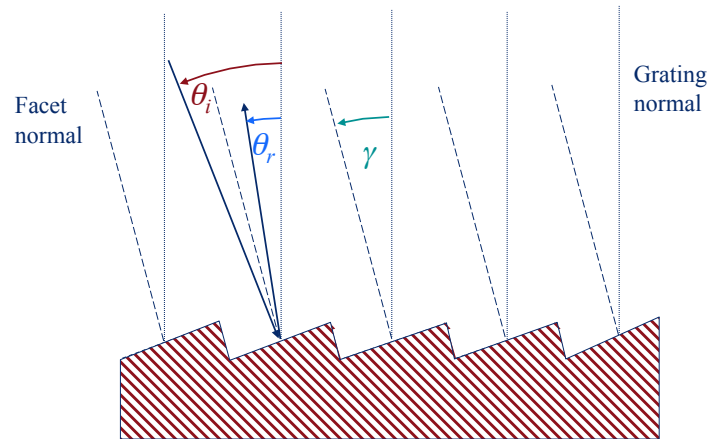


Blazed gratings

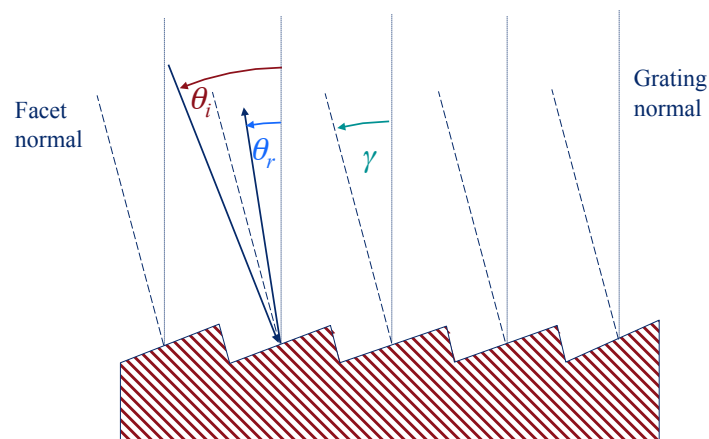
- **Shifts zeroth order light into first order.**
 - θ_i - angle of incidence of incoming beam to the grating normal
 - θ_r - angle of diffracted beam to the grating normal



- θ_i, θ_r measured from normal to grating plane.
- Peak intensity in diffraction pattern corresponds to *specular reflection off the groove facets*.
(for very polished surface, single outgoing angle)
– then $\theta_r = \theta_i = \gamma$ – the *blaze angle*.



- **Blaze angle γ related to θ_i, θ_r by $\theta_i + \theta_r = 2\gamma$ for specular reflection - the diffraction peak.**
– When $\theta_i = \theta_r = \gamma$ and $m = 1$, grating equation defines *blaze wavelength*:
 $2a \sin \gamma = \lambda_b$



linear dispersion

- **FAQ: what range of wavelengths will fit on to the detector at the spectrometer's focal plane?**

- θ_r is the same for all λ (parallel, collimated, incoming beam).

- For given order m , can differentiate the grating equation w.r.t. λ :

$$m = a \cos \theta_r \frac{d\theta_r}{d\lambda}; \quad \frac{d\theta_r}{d\lambda} = \frac{m}{a \cos \theta_r}$$

- *linear dispersion* of the spectrum in focal plane (x-direction) of the camera is then:

$$\frac{dx}{d\lambda} = f_{\text{camera}} \frac{d\theta_r}{d\lambda} = \frac{f_{\text{camera}} m}{a \cos \theta_r}$$

- e.g. Leslie Rose spectrograph

- » $a = 1/1200$ mm (1200 lines per mm), $\frac{dx}{d\lambda} = \frac{208 \times 1}{1/1200 \times \cos \theta_r} \approx \frac{249600}{\cos \theta_r}$
in first order $m = 1$, and $f_{\text{camera}} = 208$ mm, the linear dispersion is