

Lecture 5

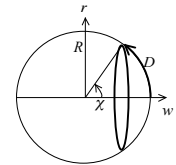
Dynamics of the Universe

$R(t) = ?$

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Robertson-Walker metric uniformly curved, evolving spacetime

$$\begin{aligned}
 ds^2 &= -c^2 dt^2 + R^2(t) (d\chi^2 + S_k^2(\chi) d\psi^2) \\
 &= -c^2 dt^2 + R^2(t) \left(\frac{du^2}{1-ku^2} + u^2 d\psi^2 \right) \\
 &= -c^2 dt^2 + a^2(t) \left(\frac{dr^2}{1-k(r/R_0)^2} + r^2 d\psi^2 \right)
 \end{aligned}$$



$$S_k(\chi) = \begin{cases} \sin \chi & (k = +1) \text{ closed} \\ \chi & (k = 0) \text{ flat} \\ \sinh \chi & (k = -1) \text{ open} \end{cases} \quad \begin{aligned} d\psi^2 &= d\theta^2 + \sin^2 \theta d\phi^2 \\ a(t) &= R(t)/R_0 \\ R_0 &= R(t_0) \end{aligned}$$

radial distance = $D(t) = R(t)\chi$
circumference = $2\pi r(t)$ $r(t) = a(t) r = R(t) u = R(t) S_k(\chi)$

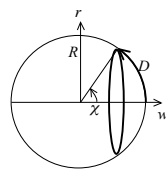
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Time and Distance vs Redshift

- We observe the **redshift** : $z = \frac{\lambda - \lambda_0}{\lambda_0} = \frac{\lambda}{\lambda_0} - 1$ λ = observed, λ_0 = emitted (rest)
- Hence we know the **expansion factor**:

$$x = 1 + z = \frac{\lambda}{\lambda_0} = \frac{\lambda(t_0)}{\lambda(t)} = \frac{R(t_0)}{R(t)} = \frac{R_0}{R(t)}$$
- When was the light emitted? $t(z) = ?$
- How far away was the source? $\chi(z) = ?$

$$\begin{aligned}
 D(t, \chi) &= R(t)\chi & D_A &= r_0(\chi)/(1+z) \\
 r(t, \chi) &= R(t)S_k(\chi) & D_L &= r_0(\chi)(1+z)
 \end{aligned}$$



- How do these depend on cosmological parameters?
 $H_0 \quad \Omega_M \quad \Omega_\Lambda$

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Time -- Redshift relation

$$\begin{aligned}
 x = 1 + z &= \frac{R_0}{R} \\
 \frac{dx}{dt} &= -\frac{R_0}{R^2} \frac{dR}{dt} \\
 &= -\frac{R_0}{R} \frac{\dot{R}}{R} \\
 &= -x H(x)
 \end{aligned}$$

Memorise this derivation!

Hubble parameter : $H \equiv \frac{\dot{R}}{R}$

$$\therefore dt = \frac{-dx}{x H(x)} = \frac{-dz}{(1+z) H(z)}$$

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Time and Distance vs Redshift

$$\frac{d}{dt} \left(x = 1 + z = \frac{R_0}{R} \right) \rightarrow dt = \frac{-dx}{x H(x)}$$

Look - back time :

$$t(z) = \int_t^{t_0} dt = \int_{1+z}^1 \frac{-dx}{x H(x)} = \int_1^{1+z} \frac{dx}{x H(x)}$$

Age : $t_0 = t(z \rightarrow \infty)$

Distance : $D = R\chi \quad r = R S_k(\chi)$

$$\chi(z) = \int d\chi = \int_t^{t_0} \frac{c dt}{R(t)} = \frac{c}{R_0} \int_1^{1+z} \frac{R_0}{R(t)} \frac{dx}{x H(x)} = \frac{c}{R_0} \int_1^{1+z} \frac{dx}{H(x)}$$

Horizon : $\chi_H = \chi(z \rightarrow \infty)$

Need to know $R(t)$, or R_0 and $H(x)$.

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Einstein's General Relativity

- 1. Spacetime geometry tells matter how to move**

- gravity = effect of **curved spacetime**
- free particles follow **geodesic** trajectories
 - $ds^2 < 0$ $v < c$ time-like massive particles
 - $ds^2 = 0$ $v = c$ null massless particles (photons)
 - $ds^2 > 0$ $v > c$ space-like tachyons (not observed)

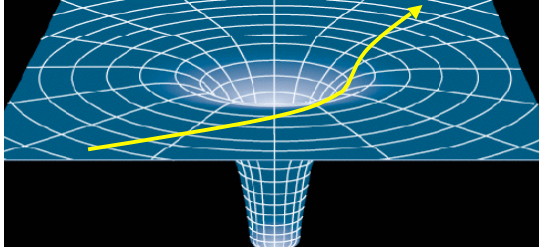
- 2. Matter (+energy) tells spacetime how to curve**

- Einstein field equations
 - nonlinear
 - second-order derivatives of metric with respect to space/time coordinates

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Geodesics

Gravity = curvature of space-time by matter/energy.
 Freely-falling bodies follow **geodesic trajectories**.
 Shortest possible path in curved space-time.
 Local curvature replaces forces acting at distance.



Space-Time Geodesics

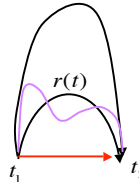
Schwarzschild Metric:

(curved space-time outside a compact mass M)

$$ds^2 = -c^2 d\tau^2 = -\left(1 - \frac{r_s}{r}\right) c^2 dt^2 + \frac{dr^2}{\left(1 - \frac{r_s}{r}\right)} + r^2 d\psi^2$$

Schwarzschild Radius: $r_s = \frac{2GM}{c^2}$

Freely-falling test particles follow a **geodesic path**.



$$\tau = \int d\tau = \int_{t_1}^{t_2} \frac{d\tau}{dt} dt$$

$$d\tau^2 = \frac{-ds^2}{c^2} = \left(1 - \frac{r_s}{r}\right) dt^2 - \frac{1}{c^2} \left[\frac{dr^2}{\left(1 - \frac{r_s}{r}\right)} + r^2 d\psi^2 \right]$$

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Space-Time Geodesics

Schwarzschild Metric:

$$ds^2 = -c^2 d\tau^2 = -\left[1 - \left(\frac{r_s}{r}\right)\right] c^2 dt^2 + \frac{dr^2}{\left[1 - \left(\frac{r_s}{r}\right)\right]} \quad r_s = \frac{2GM}{c^2}$$

Proper time:

High up clock runs faster :)

$$\frac{d\tau}{dt} = \sqrt{\frac{-ds^2}{c^2 dt^2}} = \frac{1}{c} \sqrt{-g_{tt}} = \sqrt{1 - \frac{2GM}{c^2 r}}$$

Moving clock runs slower :(

$$\frac{d\tau}{dt} = \frac{1}{c} \sqrt{-g_{tt} - g_{rr} \left(\frac{dr}{dt}\right)^2} = \sqrt{1 - \frac{2GM}{c^2 r} - \frac{v^2}{c^2} + \dots}$$

$$m c^2 \tau = \int_{t_1}^{t_2} \left(m c^2 - \frac{GMm}{r} - \frac{mv^2}{2} + \dots \right) dt = (t_2 - t_1) m c^2 - \int L(t) dt$$

$$\max[\tau] = \min[\int L dt] \quad L = E_{kin} - E_{pot} = \text{Lagrangian}$$

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Free fall =
maximum τ
proper time.

Tests of General Relativity

- **Matter bends light rays (e.g. solar eclipses)**
 - 1.75 arcsec for Sun grazing ray
- **Perihelion shift of planetary orbits**
 - GR observed
 - Mercury 43.03 43.11 +/- 0.45 arcsec/century
 - Venus 8.6 8.4 +/- 4.8
 - Earth 3.8 5.0 +/- 1.2
 - Icarus 10.3 9.8 +/- 0.8
- **Time delay of Venus radar reflections**
 - ~200 microsec
- **Decay of orbit of binary pulsars**
 - due to emission of gravitational waves
- **All tests on small size scales.**
- **For cosmology, assume GR valid on larger scales.**

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