

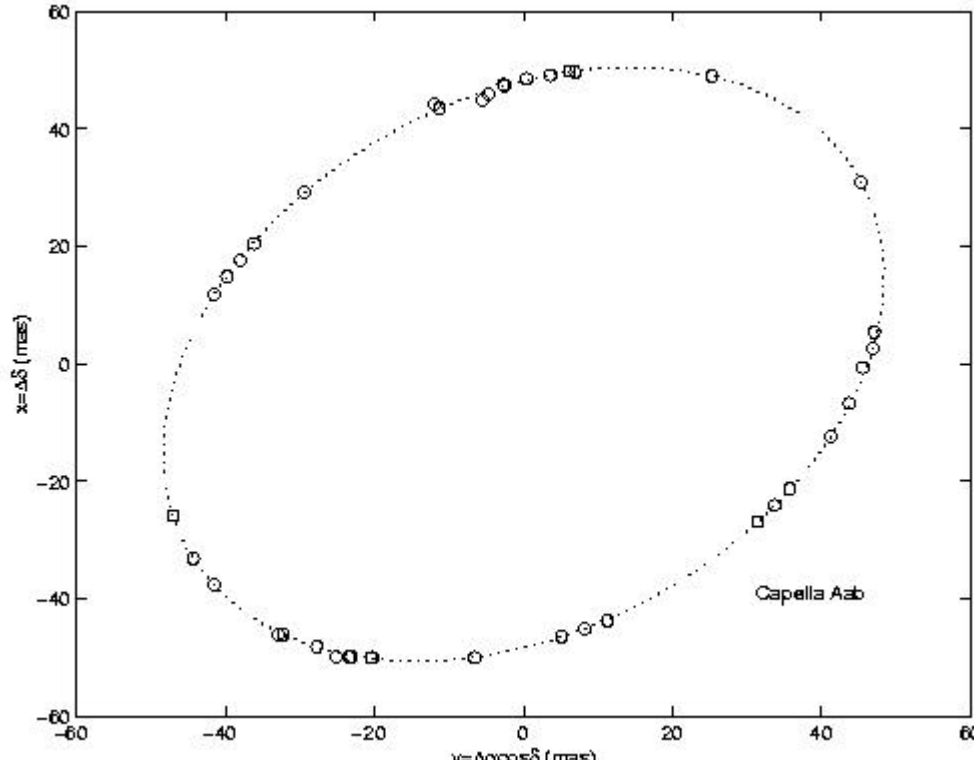
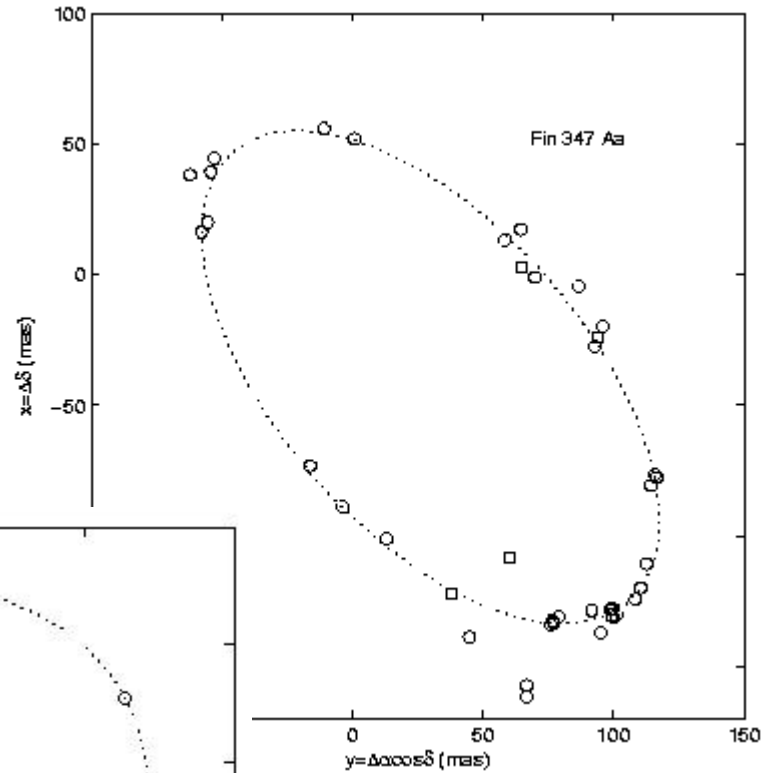
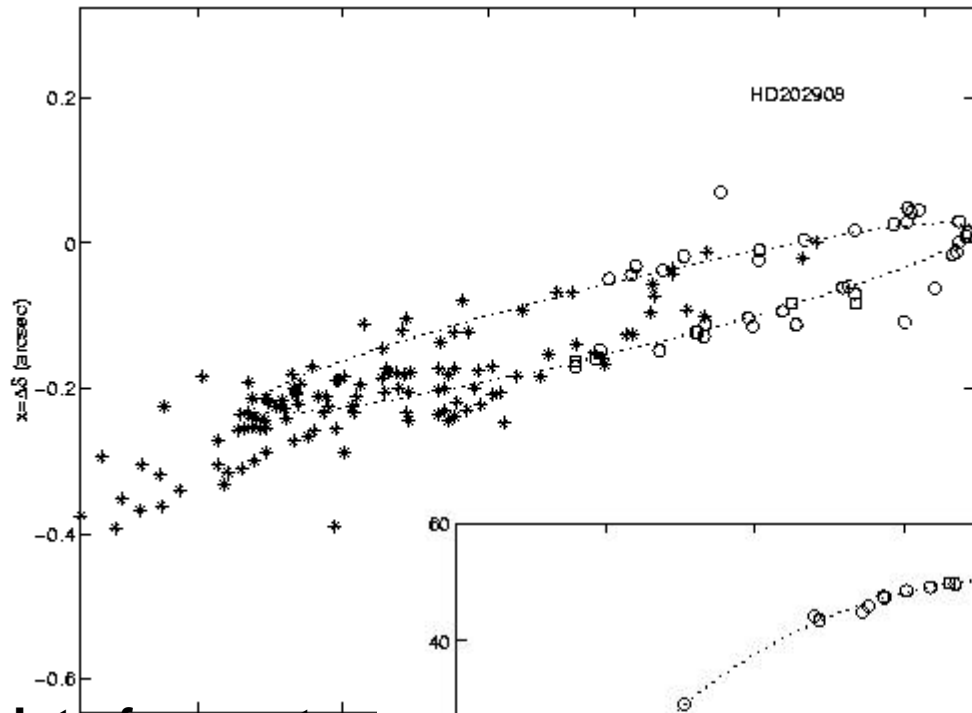
# Inclinations

- **Radial velocities give minimum masses**

$$\text{SB2: } m \sin^3 i \quad \text{SB1: } \frac{m^3 \sin^3 i}{M^2}$$

- **need inclinations**
  - astrometric orbits
  - eclipsing binaries
  - polarimetry
  - apsidal motion

# Astrometric Orbits



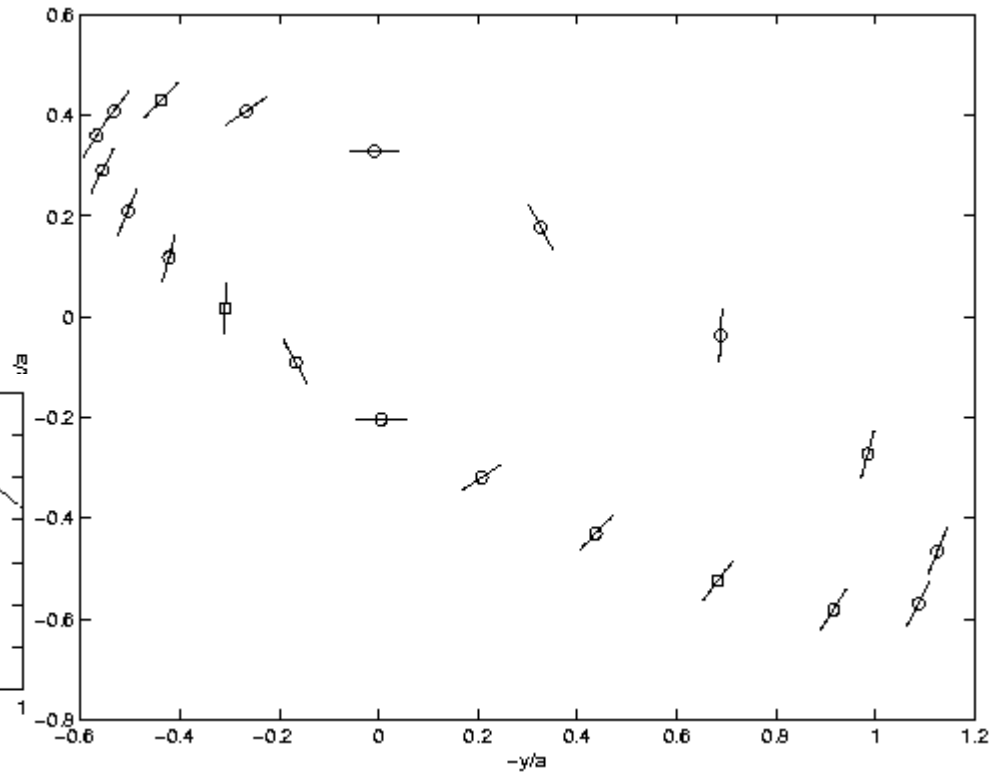
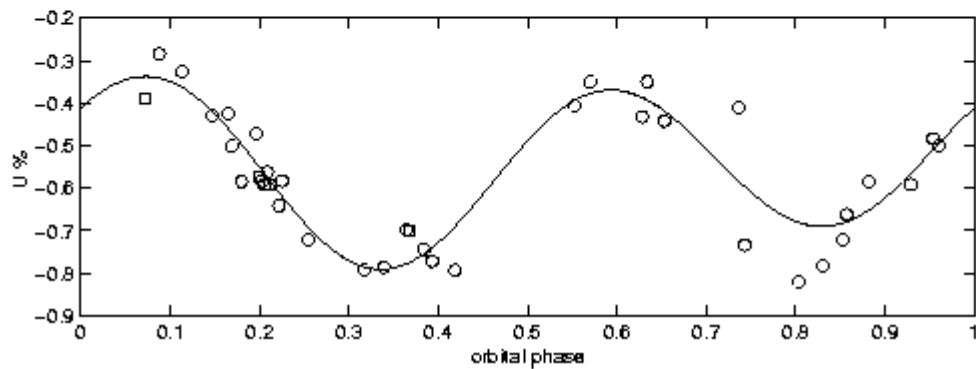
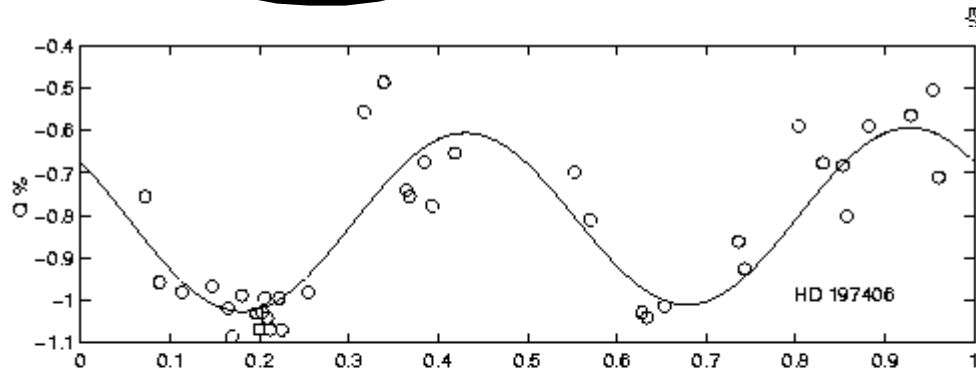
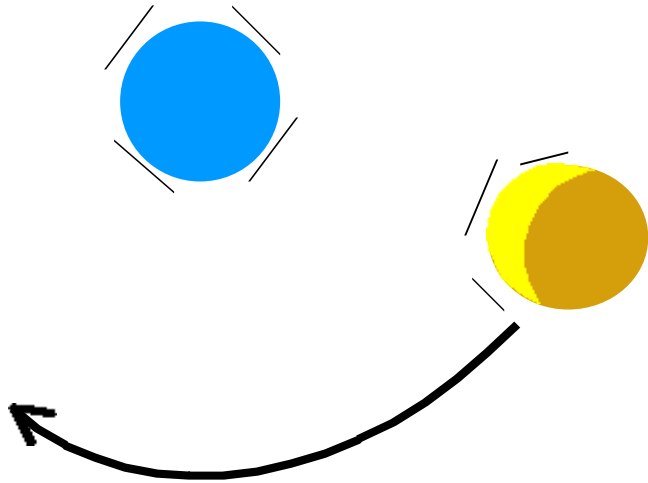
**Interferrometry**  
measures the  
relative orbit  
(milli-arcsec  
accuracy).

**Fit elliptical  
orbit model to  
 $x(t)$  and  $y(t)$ .**

**Get inclination,  
hence masses**

# Polarimetry

Scattered light is polarised (e.g. our sky)

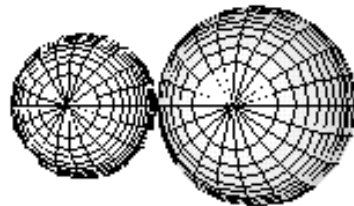


Get inclination,  
hence masses

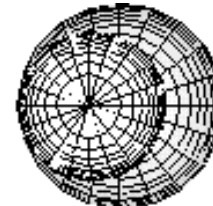
# Eclipsing Binaries

- **Binary properties from eclipses**
  - sizes, shapes of stars, inclination,
  - temperatures, limb darkening, apsidal motion
  - 4 contact phases:

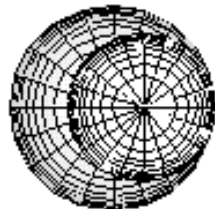
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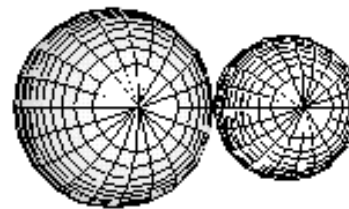
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3

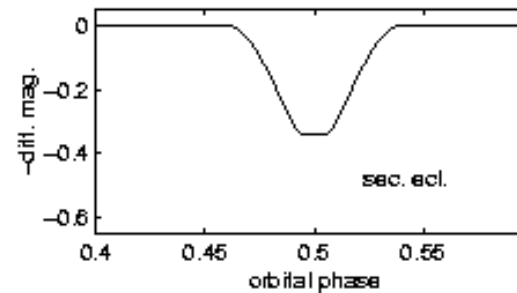
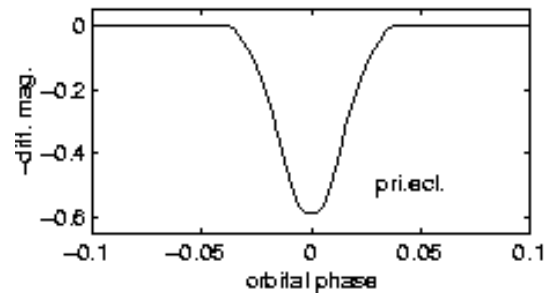


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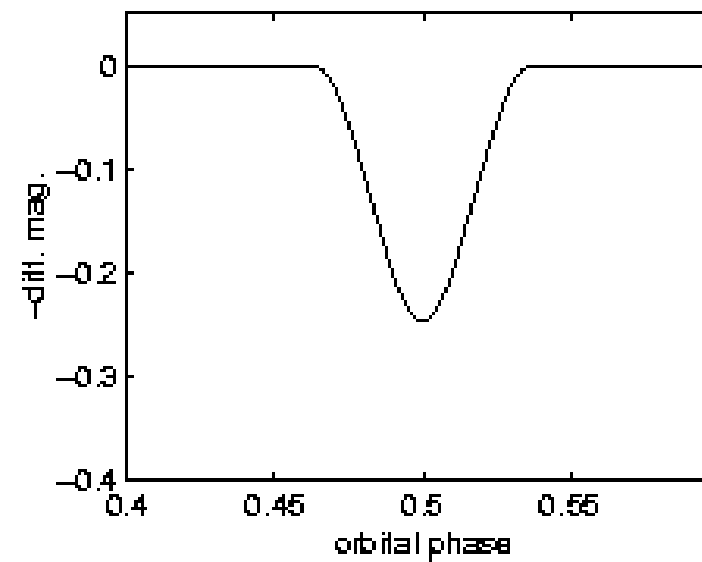
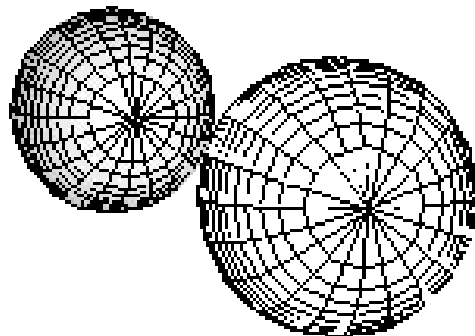
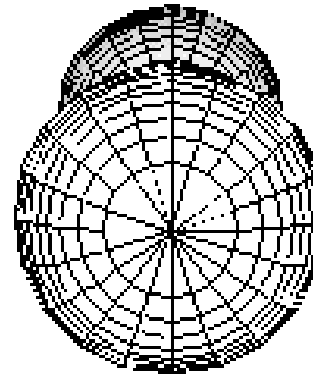
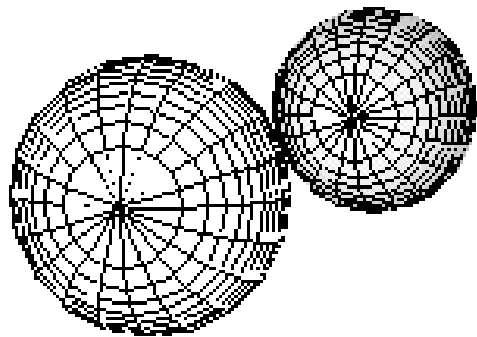


primary eclipse

secondary eclipse



# Partial Eclipse



# Eclipses

$$R_1 > R_2 \quad \mathbf{q} \equiv 2p \mathbf{f}$$

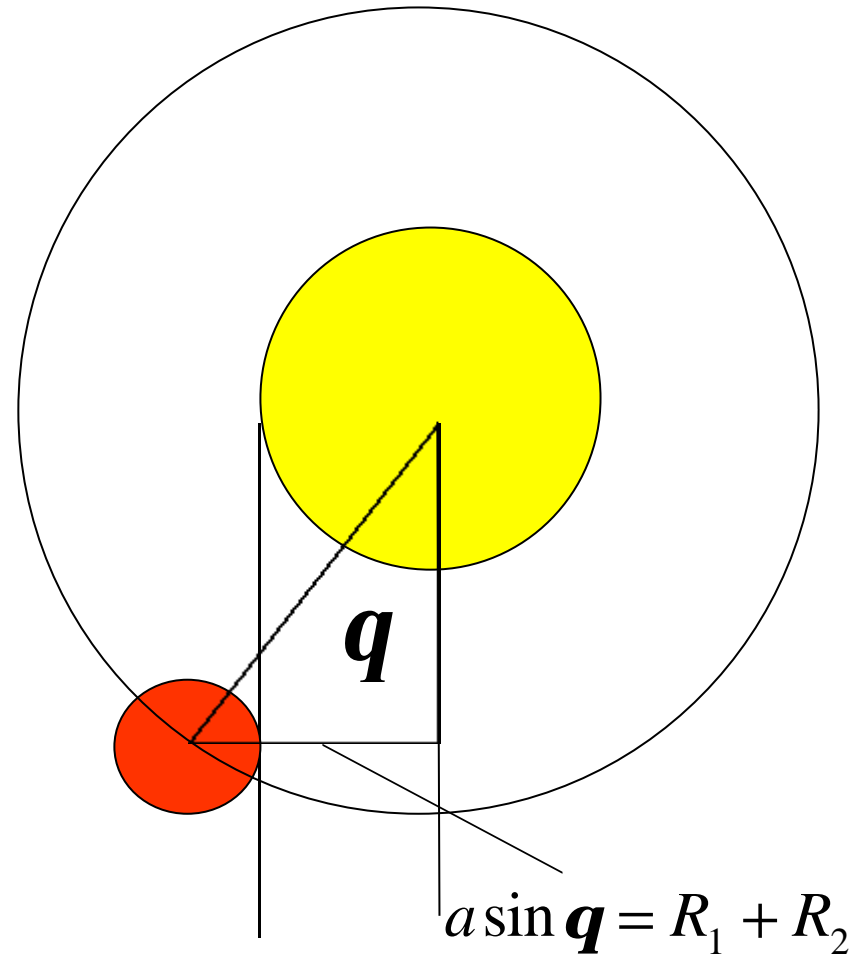
edge-on ( $i = 90^\circ$ ):

4 contact phases:

$$R_1 \pm R_2 = a \sin \mathbf{q}$$

$$\frac{R_2}{a} = \frac{\sin \mathbf{q}_1 - \sin \mathbf{q}_2}{2}$$

$$\frac{R_1}{a} = \frac{\sin \mathbf{q}_1 + \sin \mathbf{q}_2}{2}$$



# Eclipses

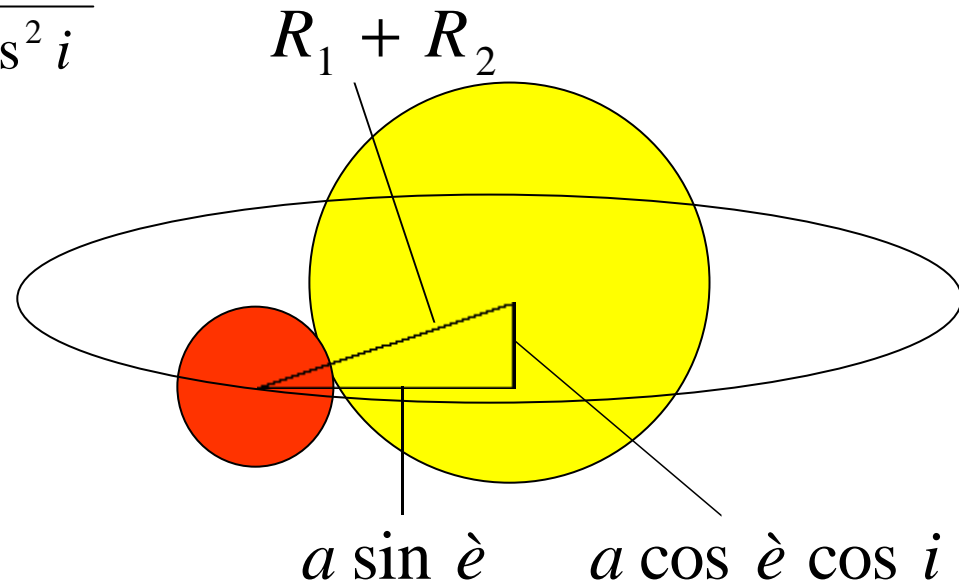
4 contact phases :

$$R_1 \pm R_2 = a \sqrt{\sin^2 \mathbf{q} + \cos^2 \mathbf{q} \cos^2 i}$$

$$= a \sqrt{1 + \cos^2 \mathbf{q} \sin^2 i}$$

4 measurements :  $\mathbf{f}_1$   $\mathbf{f}_2$   $\mathbf{f}_3$   $\mathbf{f}_4$

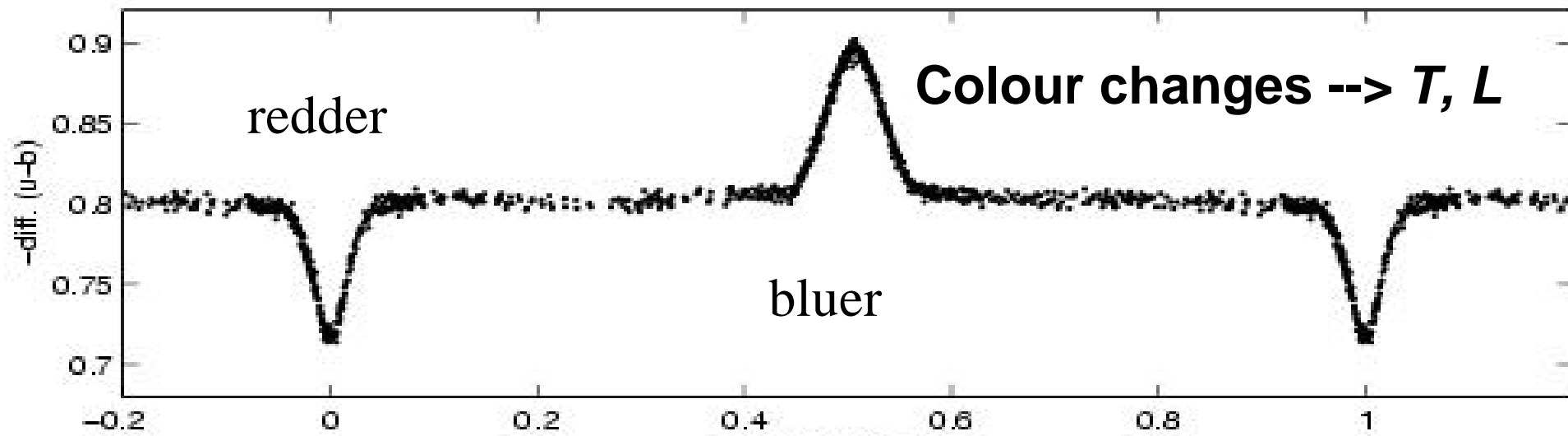
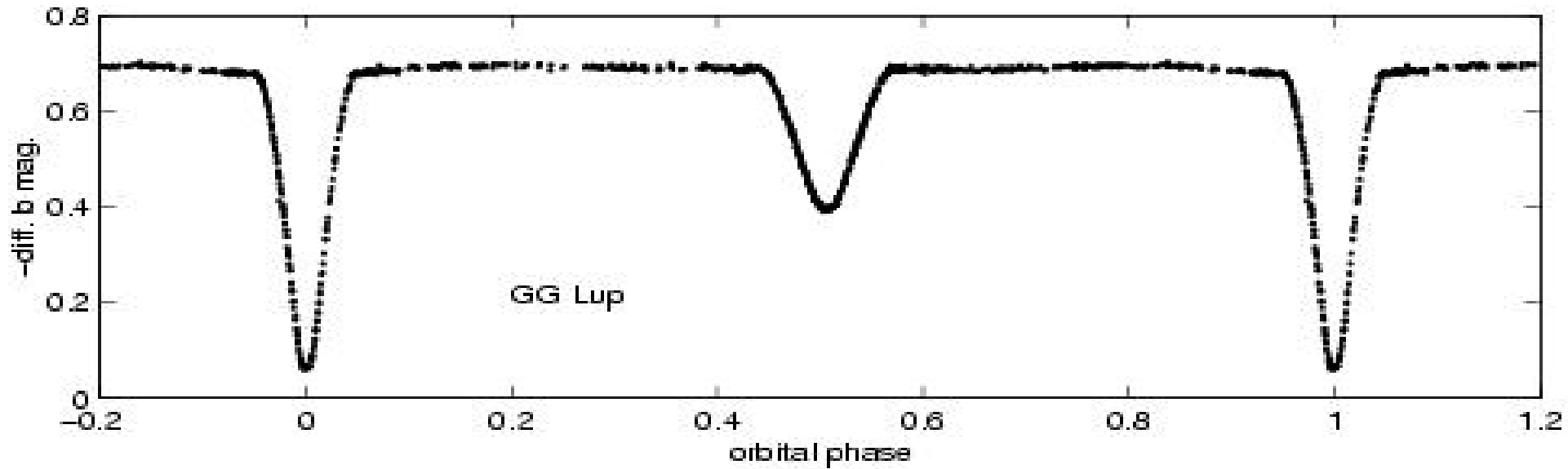
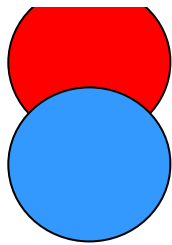
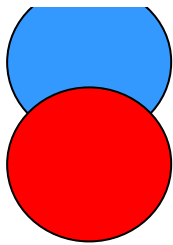
4 parameters :  $\frac{R_1}{a}$   $\frac{R_2}{a}$   $i$   $\mathbf{f}_0$



mid eclipse :  $\mathbf{q} = 0$

total eclipse :  $a \cos i < R_1 - R_2$

partial eclipse :  $R_1 - R_2 < a \cos i < R_1 + R_2$



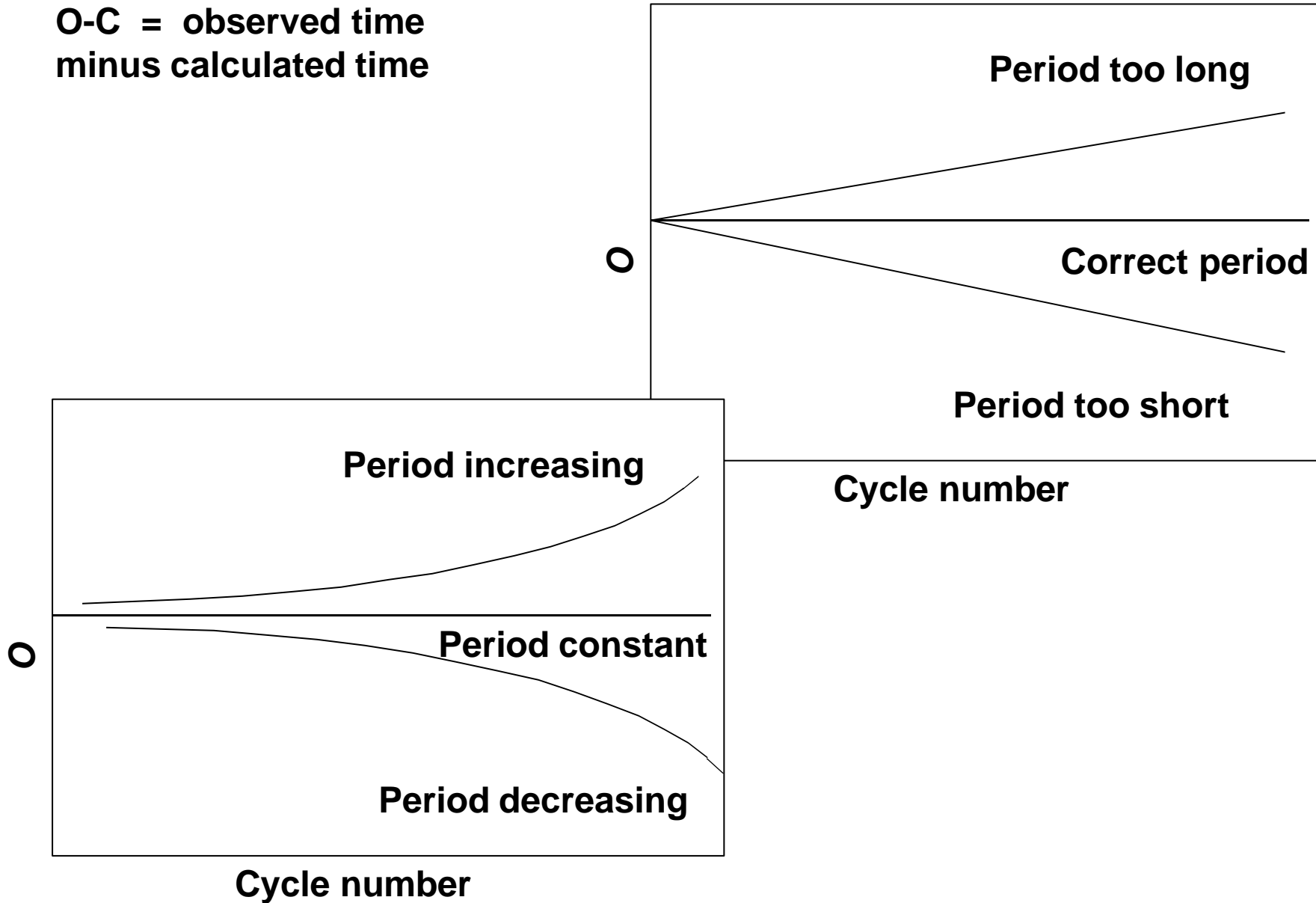


# Application to Binary Pulsars

- **binary system where one star is a pulsar**
  - emits ‘pulses’ of radiation
  - accurate timing possible (accurate clocks)
  - need narrow pulses
    - radio signals from neutron stars
- **solitary pulsar**
  - if at 0 velocity relative to us
    - time between pulses,  $dt = \text{constant}$  (unless being spun up/down)
  - if at  $V_{\text{rel}}$  relative velocity
    - $dt = \text{constant} \times \text{pulse number}$
  - if pulsar spins up,  $dt$  decreases with pulse number
    - concave curve
  - if pulsar spins down,  $dt$  increases with pulse number
    - convex curve

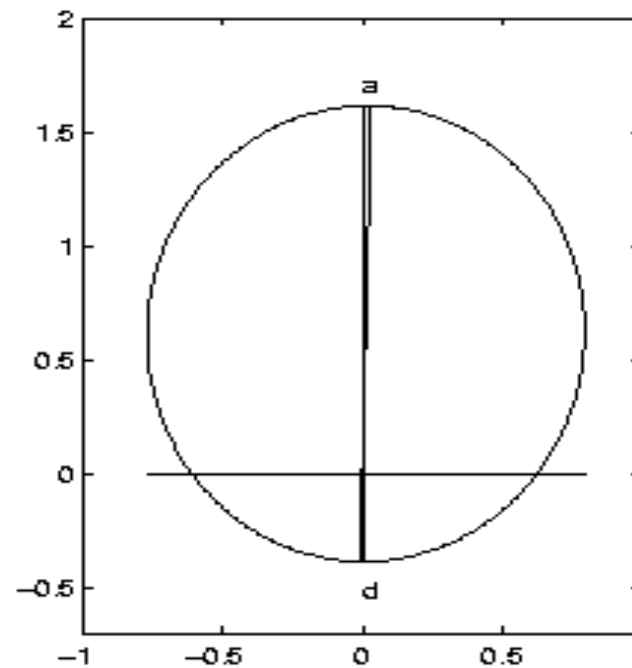
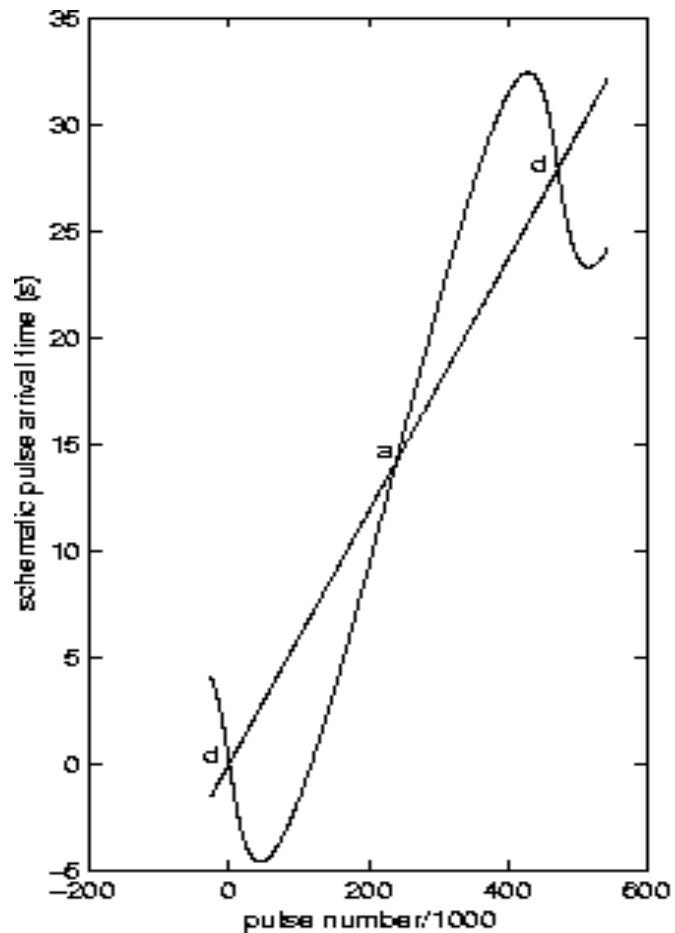
# Timing Residuals

O-C = observed time  
minus calculated time



# Binary Pulsars

- In binary system, time between pulses affected by orbital motions
  - due to light travel time (distance) changing along orbit



# Light travel time

pulsar orbit :  $r_p = \frac{a_p (1 + e^2)}{1 + e \cos \mathbf{q}}$

distance along line of sight :

$$z_p = r_p \sin i \sin(\mathbf{q} + \mathbf{w})$$

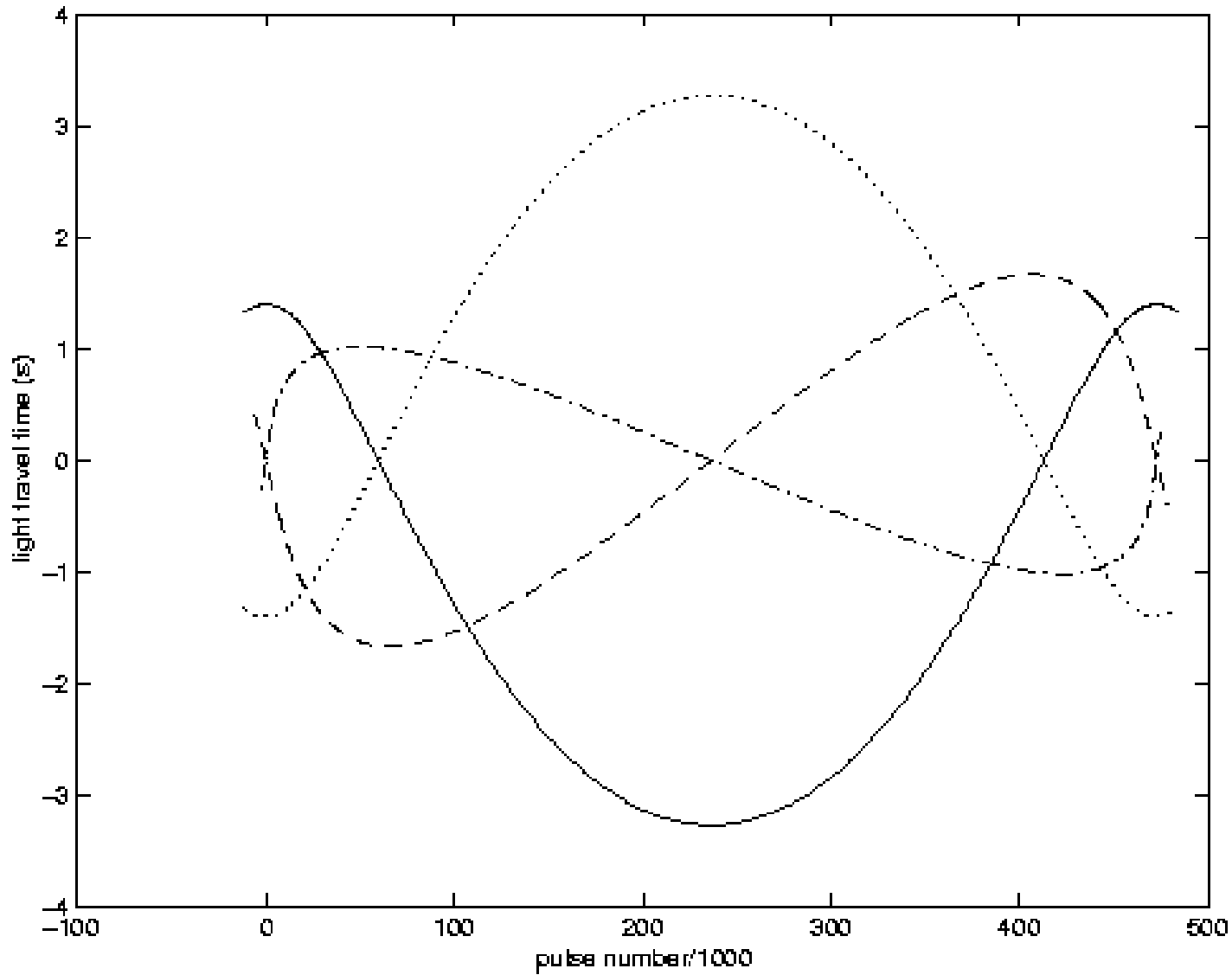
light travel time :

$$\frac{z_p}{c} = \frac{a_p \sin i}{c} (1 + e^2) \frac{\sin(\mathbf{q} + \mathbf{w})}{1 + e \cos \mathbf{q}}$$

circular orbit

$$\frac{z_p}{c} = \frac{a_p \sin i}{c} \sin\left(\frac{2\mathbf{p}}{P} (t - T_0)\right)$$

# light travel time



# Binary Pulsar timing residuals

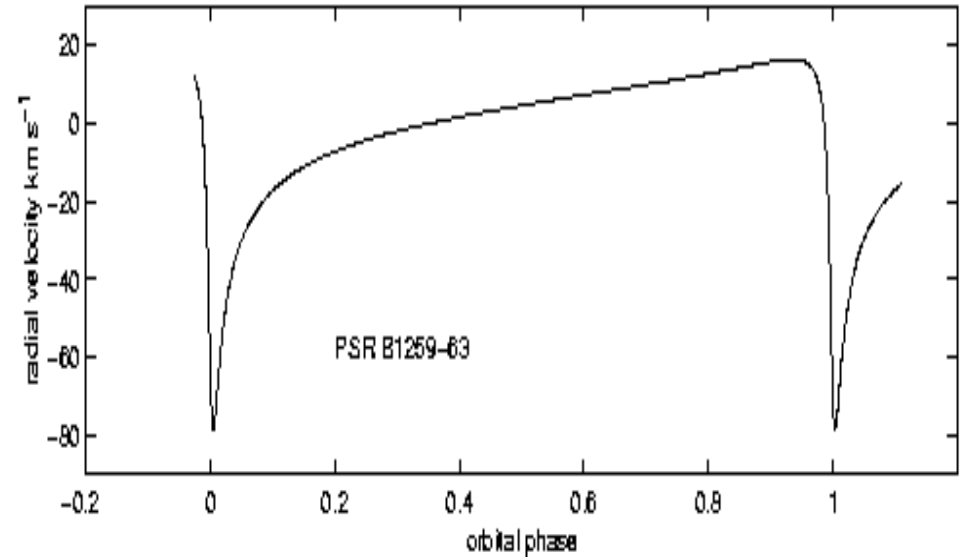
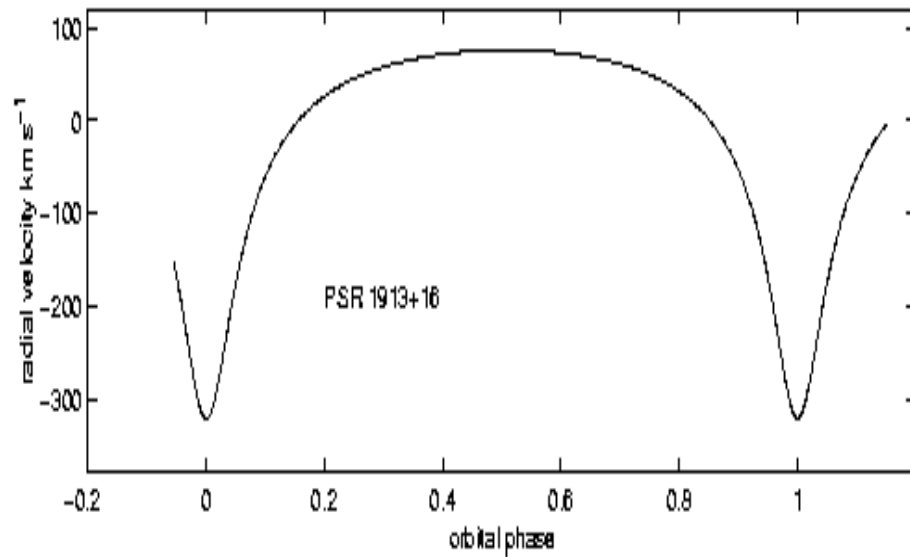
- Time difference between predicted ,  $j_n \mathbf{t}$ , and actual (binary) pulse arrival times,  $t_n$  is

$$\Delta t = t_n - j_n \mathbf{t} = a t + b \sin \left[ \frac{2p (t - T_0)}{P} \right]$$

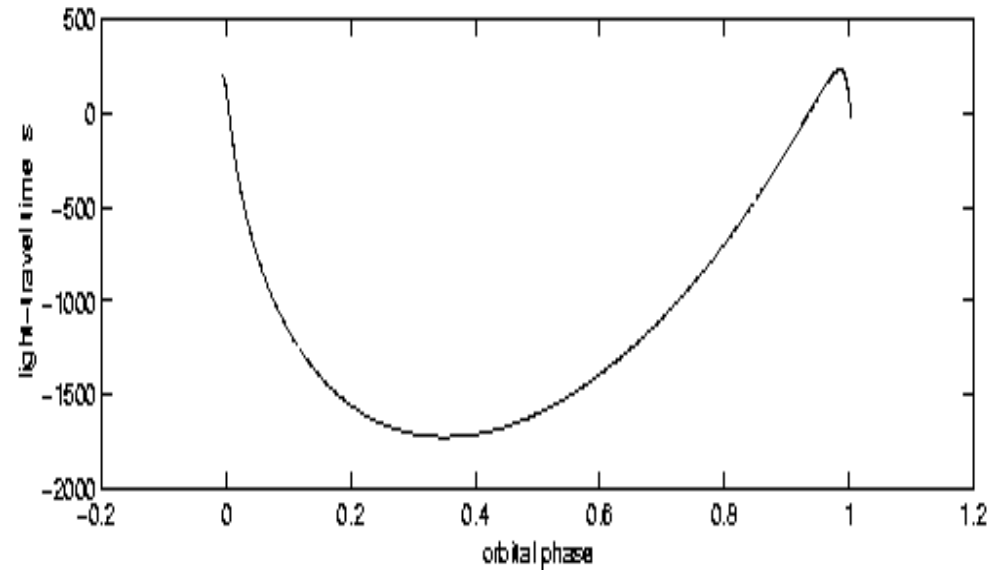
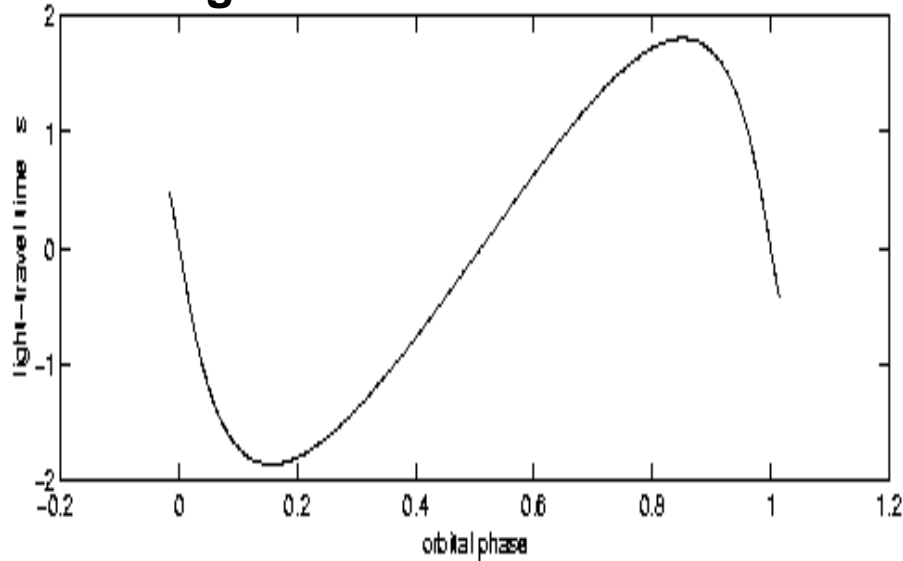
- $P$  is the orbital period,  $T_0$  is a reference time
- $a, b$  are determined by the velocity of the pulsar
  - $a$ : from systematic velocity
  - $b$ : from Keplerian velocity
- for circular orbits:  $b = (a_p/c) \sin i$

# Binary Pulsar Orbits

radial velocity



timing



# Mass determinations

- **visible companion star**

- O-B star in High-Mass X-ray Binaries (HMXB)
- A-K star in Low-Mass X-ray Binaries (LMXB)

$$a_c \sin i = \frac{(1-e^2)^{1/2} K_c P}{2p}$$

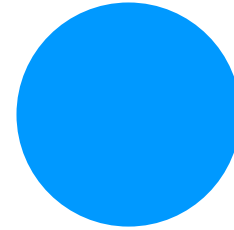
mass function

$$f(m_p) = \frac{m_p^3 \sin^3 i}{M^2} = \frac{(1-e^2)^{3/2} K_c^3 P}{2p G}$$

mass ratio,  $q$ ,

$$q = \frac{m_p}{m_c} = \frac{a_c \sin i}{a_p \sin i}$$

- If inclination,  $i$ , can be found, then masses follow





# Frequency shifts

- **Binary orbit also affects pulsar frequency**
  - radio pulsars , very narrow pulse widths
  - pulse frequency affected by orbital velocity
  - Doppler shift:

$$\Delta f = f \frac{V_{rad}}{c} = f \frac{\dot{z}}{c}$$

- gives a phase lag of:

$$\begin{aligned} \Delta \mathbf{f} &= \int_{T_0}^t \Delta f dt \approx f_0 \frac{\dot{z}}{c} (t - T_0) \\ &= f_0 \frac{\dot{z}}{c} \left[ \frac{z}{c} \right] \end{aligned}$$

# Pulsar Phase lag

- **Combined phase lag is**
  - from light travel time due to orbit

$$\Delta f_L = -f \frac{z}{c}$$

- and from Doppler shift

$$\Delta f_D = f_0 \frac{\dot{z}}{c} \left[ \frac{z}{c} \right]$$

- hence

$$\Delta f = \Delta f_D + \Delta f_L \approx \left[ \frac{z}{c} \right] \left[ f_0 \frac{\dot{z}}{c} - f_0 - \dot{f}_0 (t - T_0) \right]$$

- generally  $\Delta f_D \cong 0.001 \Delta f_L$

- but measurable in radio pulsars