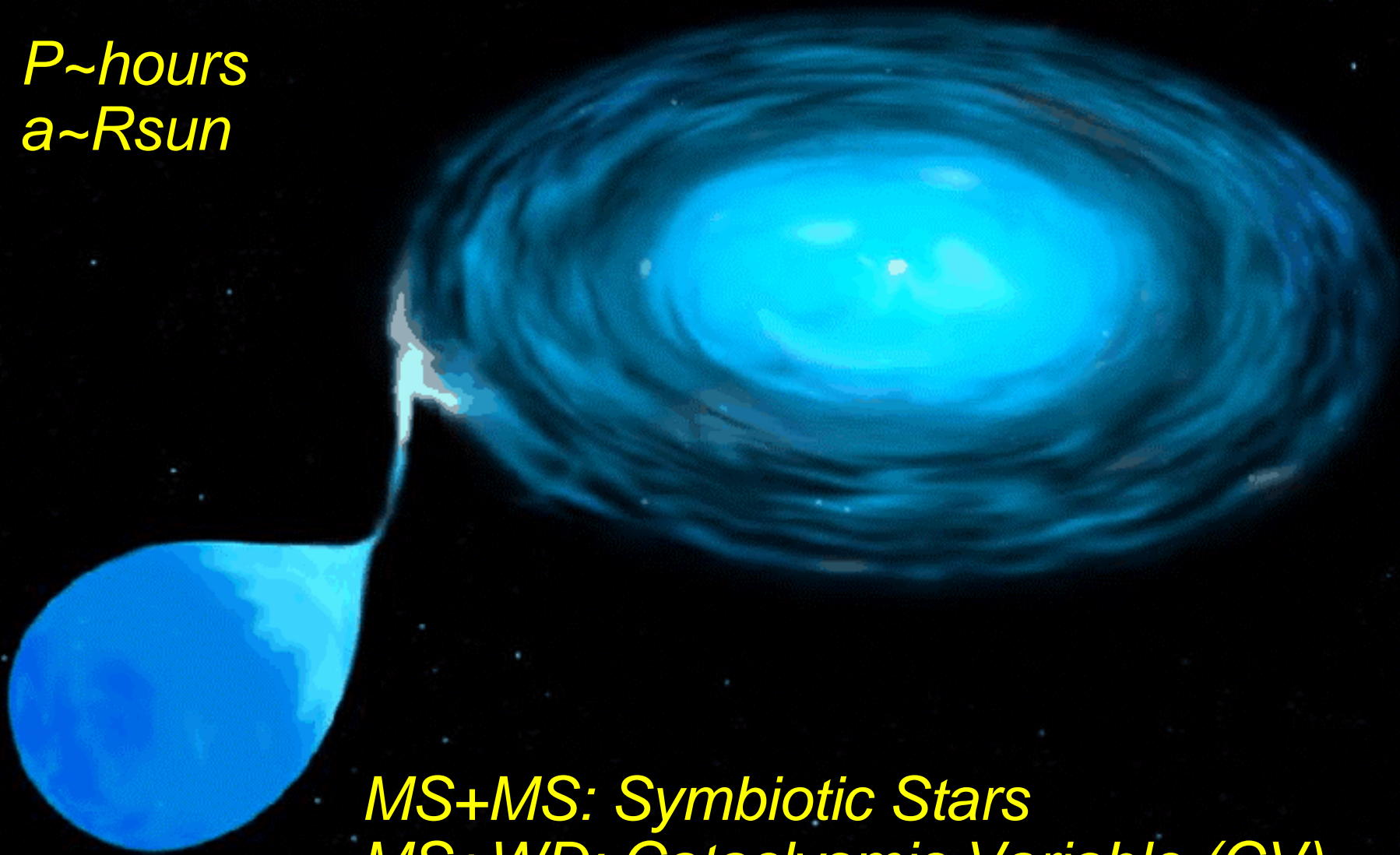


Binary Stars with Accretion Discs

P ~ hours
a ~ R_{sun}



MS+MS: Symbiotic Stars
MS+WD: Cataclysmic Variable (CV)
MS+NS or BH: Low-Mass X-ray Binary (LMXB)

Importance of Accretion Discs

- **Formation of compact objects**
 - friction moves angular momentum outward
 - allows matter to spiral inward
 - build up compact object at centre

- **Generation of light**
 - gravitational potential energy
 - converted by friction to heat
 - radiated as light

3 Classic Papers

- **Black hole accretion discs**

Shakura, Sunyaev 1973 A&A 24 337

- **Time-dependent discs**

Pringle, Lynden-Bell 1974 MNRAS 168 603

- **Gas streams**

Lubow, Shu 1975 ApJ 198 383

Accretion Energy

- **Nuclear energy (H -> He) Efficiency:**

$$\frac{\Delta E_{nuc}}{m} \approx 6 \times 10^{14} \frac{\text{J}}{\text{kg}}$$

$$h = \frac{\Delta E_{nuc}}{m c^2} = 0.7\%$$

- **Accretion energy**

$$\begin{aligned} \frac{\Delta E_{acc}}{m} &\approx \frac{G M}{R} \approx 1.7 \times 10^{11} \left(\frac{M}{M_{sun}} \right) \left(\frac{R}{R_{sun}} \right)^{-1} \frac{\text{J}}{\text{kg}} \\ &\approx 1.7 \times 10^{16} \left(\frac{M}{1.4 M_{sun}} \right) \left(\frac{R}{10 \text{ km}} \right)^{-1} \frac{\text{J}}{\text{kg}} \end{aligned}$$

- **more efficient for compact objects**

Efficiency

- Nuclear energy (H -> He) Efficiency:

$$\Delta E_{nuc} = \mathbf{h} m c^2 \quad \mathbf{h} = 0.7\%$$

- Accretion efficiency

$$L \approx \frac{G M \dot{M}}{R} = \mathbf{h} \dot{M} c^2 \quad \mathbf{h} = \frac{G M}{R c^2} = \frac{R_H}{2 R} \quad R_H \approx 3 \text{ km} \left(\frac{M}{M_{sun}} \right)$$

- Compactness --- M / R

- Black hole - smallest stable orbit

– non-rotating $R = 6 R_H \quad \mathbf{h} = \frac{1}{12} \approx 8\%$

– maximally-rotating $R = 3 R_H \quad \mathbf{h} = \frac{1}{6} \approx 17\%$

- White Dwarf

$$M \approx M_{sun} \quad R \approx 10^4 \text{ km} \quad \mathbf{h} \sim 10^{-4}$$

Radiation Opposes Gravity

- Coulomb attraction ties electrons and protons
- gravity pulls in

$$F_{grav} = -\frac{G M m}{r^2}$$

- radiation pushes electrons out

$$F_{rad} = \frac{L \mathbf{s}_T}{4\pi r^2 c} \quad L_{acc} = -\frac{G M \dot{M}}{R}$$

Thompson electron
scattering cross-section

$$\mathbf{s}_T = 6.7 \times 10^{-21} \text{ m}^2$$

- Notes:
 - radiation pressure opposes gravity
 - same R scaling
 - radiation pressure > gravity at high Mdot

Eddington Luminosity

- Photon momentum

$$p_g = \frac{h \mathbf{n}}{c}$$

- Density of photons

$$n_g = \frac{L / h \mathbf{n}}{4\pi R^2 c}$$

- Radiative force (per electron)

$$F_{rad} = n_g p_g \mathbf{s}_T c = \frac{L \mathbf{s}_T}{4\pi R^2 c} \quad \mathbf{s}_T = 6.7 \times 10^{-25} \text{ cm}^2$$

- Total force (per electron + proton pair)

$$F_{rad} + F_{grav} = \frac{L \mathbf{s}_T}{4\pi R^2 c} - \frac{G M (m_p + m_e)}{R^2}$$

- Eddington Luminosity

$$L_{Edd} = \frac{4\pi G M (m_p + m_e) c}{\mathbf{s}_T} = 1.3 \times 10^{38} \left(\frac{M}{M_{sun}} \right) \text{ erg s}^{-1}$$

Eddington Accretion Rate

- Eddington Luminosity

$$L_{Edd} = \frac{4\mathbf{p} G M m_p c}{\mathbf{s}_T}$$

- Accretion Luminosity

$$L_{acc} = \frac{G M \dot{M}}{R}$$

- Eddington Accretion Rate

$$\frac{L_{acc}}{L_{Edd}} = \frac{\mathbf{s}_T \dot{M}}{4\mathbf{p} m_p c R} = \frac{\dot{M}}{\dot{M}_{Edd}}$$

$$\dot{M}_{Edd} = \frac{4\mathbf{p} m_p c R}{\mathbf{s}_T}$$

$$\approx 10^{-5} \left(\frac{R}{R_{sun}} \right) \frac{M_{sun}}{\text{yr}}$$

	kg/s	M_{sun} / yr
WD	6×10^{19}	10^{-7}
NS	10^{12}	1.4×10^{-12}

Rough Temperatures

- **Optically thick**

- **Blackbody radiates accretion luminosity**

$$\mathbf{s} T_b^4 = \frac{L_{acc}}{4\mathbf{p} R^2} \quad T_b = \left(\frac{L_{acc}}{4\mathbf{p} R^2 \mathbf{s}} \right)^{1/4} \approx \left(\frac{G M \dot{M}}{4\mathbf{p} R^3 \mathbf{s}} \right)^{1/4}$$

- **Optically thin**

- **potential energy released**
- **= thermal energy of shocked gas**

$$\frac{G M m_p}{R} = 2 \times \frac{3}{2} k T_{th} \quad T_{th} \approx \frac{G M m_p}{3 k R}$$

- **Radiation temperature and photon energy**

$$T_b < T_{rad} < T_{th} \quad k T_b < h \mathbf{n} < k T_{th}$$

Rough Temperatures

- **Neutron Star or Black Hole**

$$L_{acc} \sim L_{Edd} \sim 10^{38} \text{ erg s}^{-1}$$

$$10^7 \text{ K} < T_{rad} < 10^{11} \text{ K}$$

$$1 \text{ keV} < h\mathbf{n} < 50 \text{ MeV}$$

- mid to hard X-rays

- **White Dwarf**

$$L_{acc} \sim 10^{33} \text{ erg s}^{-1}$$

$$6 \times 10^4 \text{ K} < T_{rad} < 10^9 \text{ K}$$

$$6 \text{ eV} < h\mathbf{n} < 100 \text{ keV}$$

- optical - uv - soft X-ray

Roche Lobe Overflow

- **Initial stream velocity**

$$V_{\parallel} \sim c_s \sim 10 \text{ km s}^{-1}$$

- **L1 velocity relative to centre of mass**

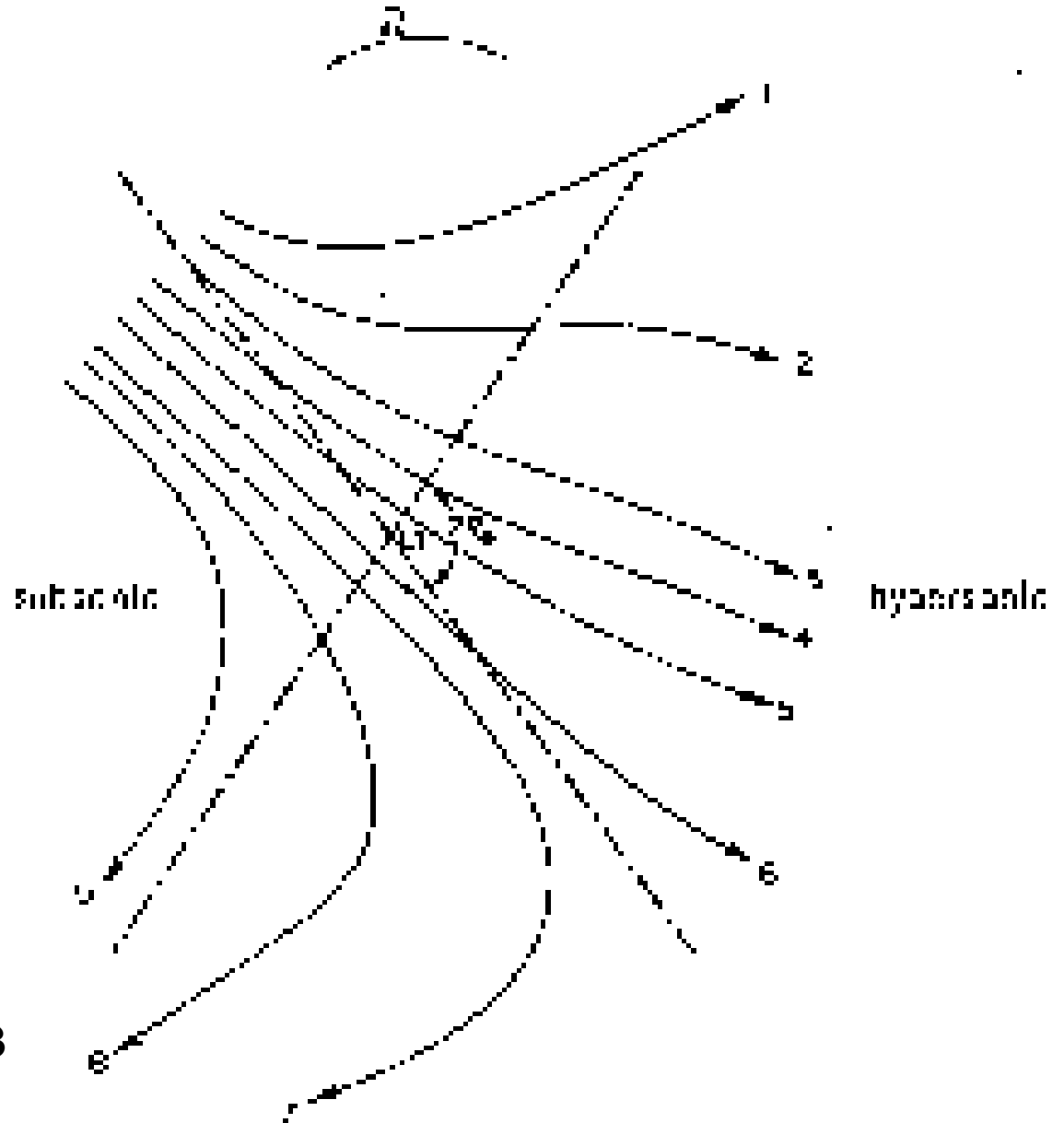
$$V_{\perp} = \frac{2p}{P} R_{L1} \approx \frac{2p}{P} a (0.5 - 0.23 \log q)$$

$$\sim 100 \text{ km s}^{-1} \left(\frac{M}{M_{sun}} \right)^{1/3} \left(\frac{P}{\text{day}} \right)^{-1/3}$$

- **subsonic --> supersonic transition at nozzle**
- **ballistic trajectory in Roche potential**
 - (neglect pressure forces)

Flow thru the L1 nozzle

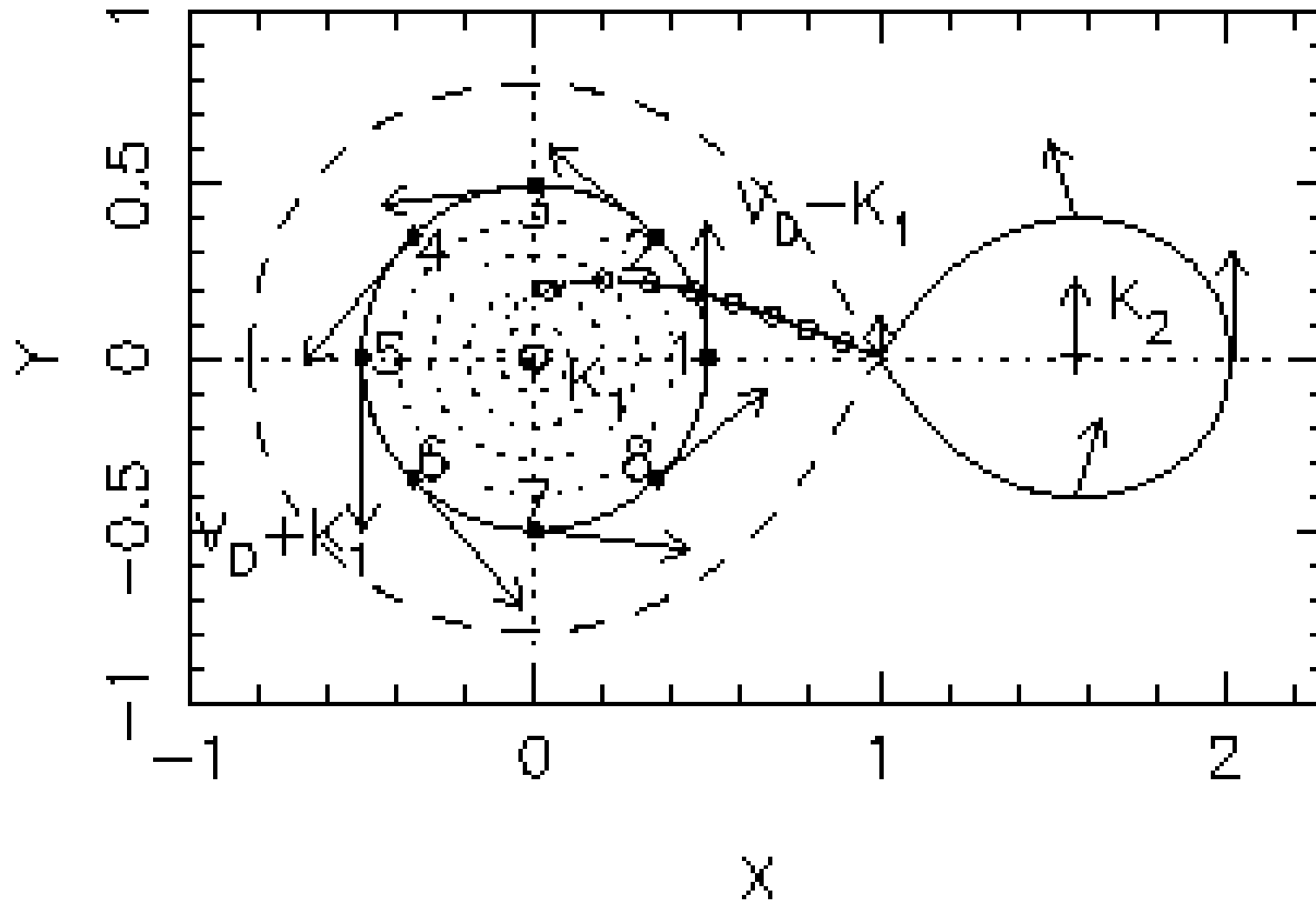
subsonic \rightarrow hypersonic



Lubow, Shu 1975 ApJ 198 383

Balistic Stream Trajectory

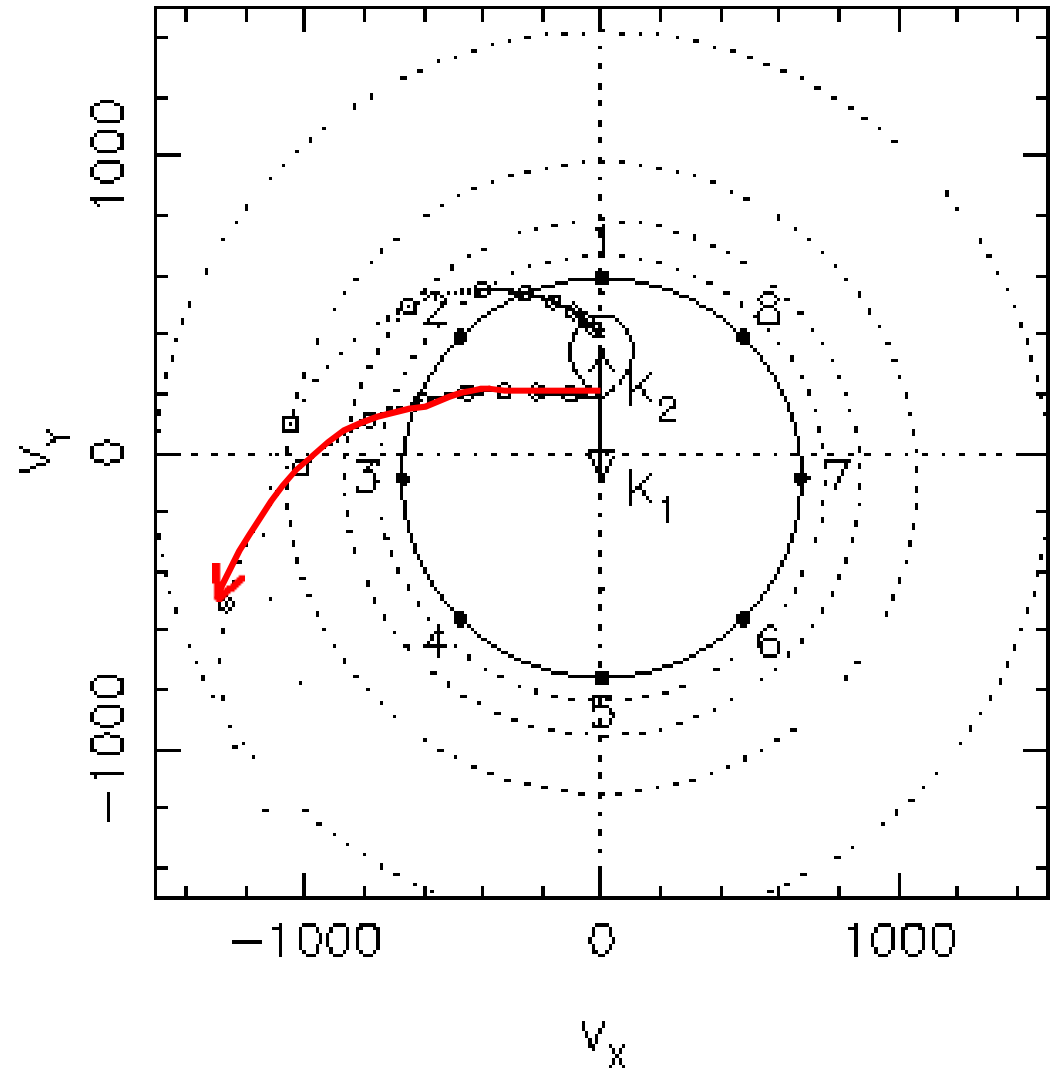
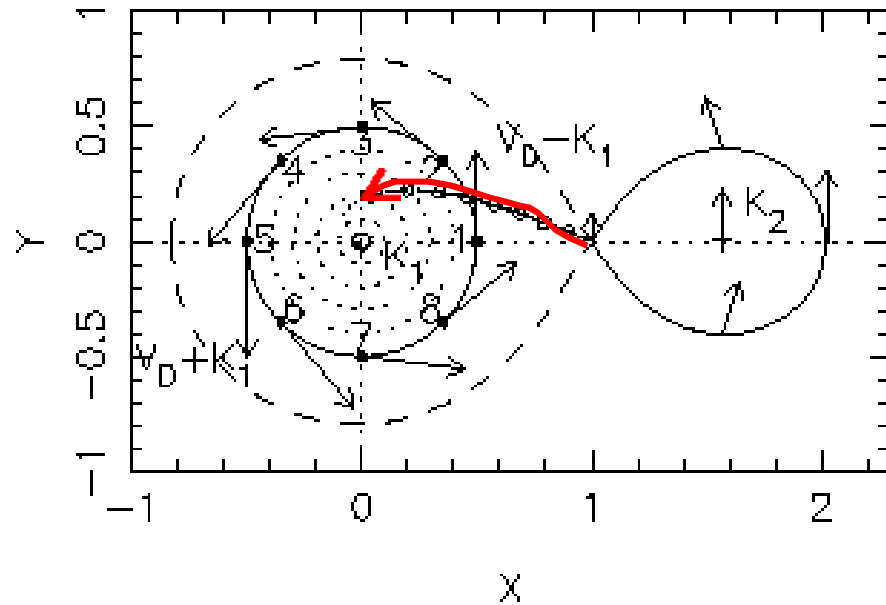
Position Coordinates



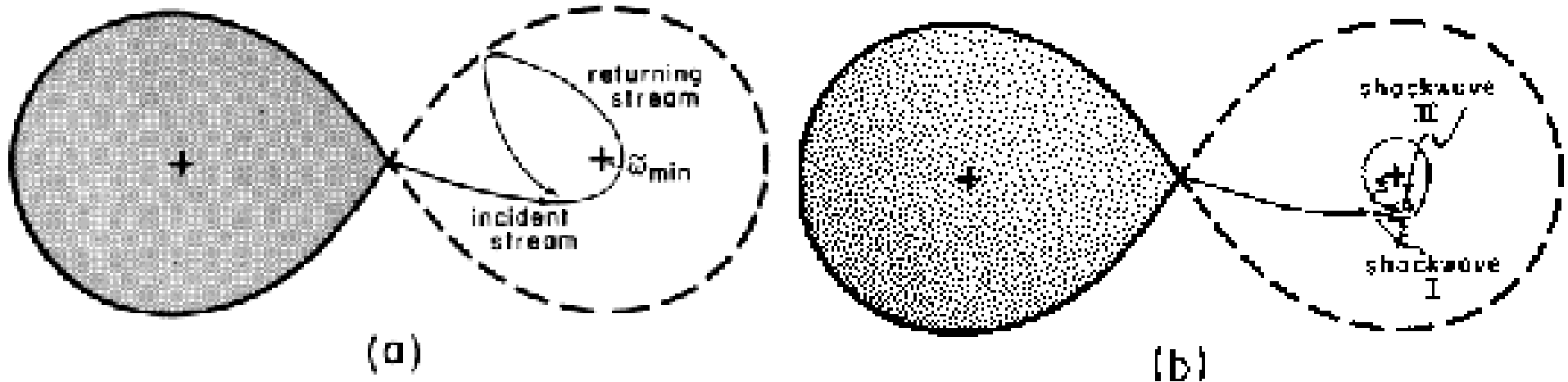
Balistic Stream Trajectory

Doppler Coordinates

Position Coordinates



Formation of a Ring

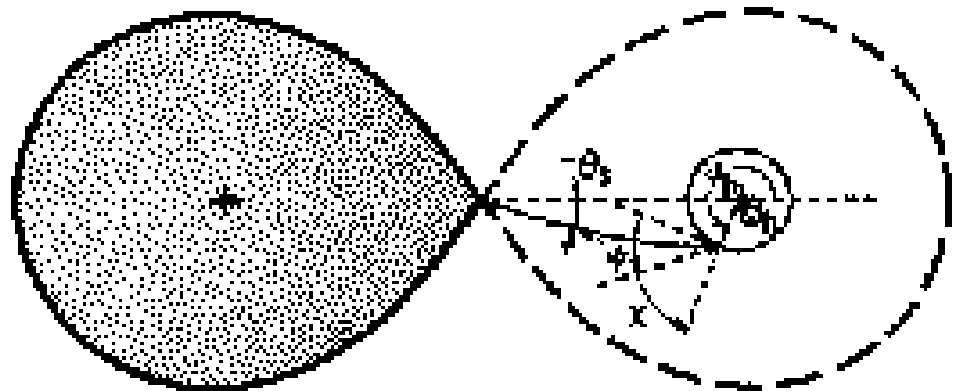


Circularisation Radius

same angular momentum as L1

$$(GM R_{circ})^{1/2} = \frac{2P}{P} R_{L1}^2$$

$$\frac{R_{circ}}{a} = (1+q) \left(\frac{R_{L1}}{a} \right)^4 \approx (1+q) (0.5 - 0.23 \log q)^4 \quad (c)$$



Friction -> Spreading

- **Kepler Velocity**

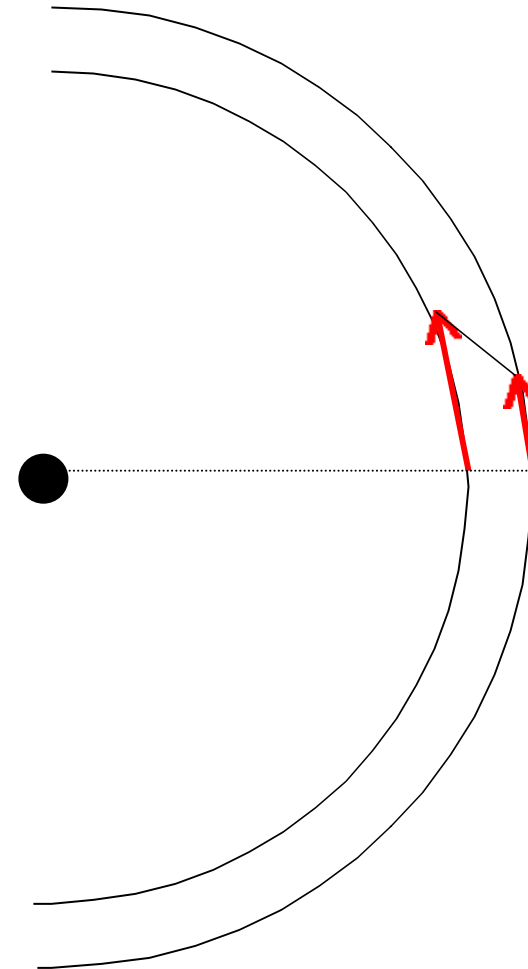
$$V = \sqrt{\frac{G M}{R}}$$

- **Differential rotation (Shear)**

$$\begin{aligned}\Delta V &= \frac{d}{dR} \left(\sqrt{\frac{G M}{R}} \right) \Delta R \\ &= \frac{V}{2} \frac{\Delta R}{R}\end{aligned}$$

- **Friction**

- opposes shear
- causes ring to spread
- inward + outward
- diffusion



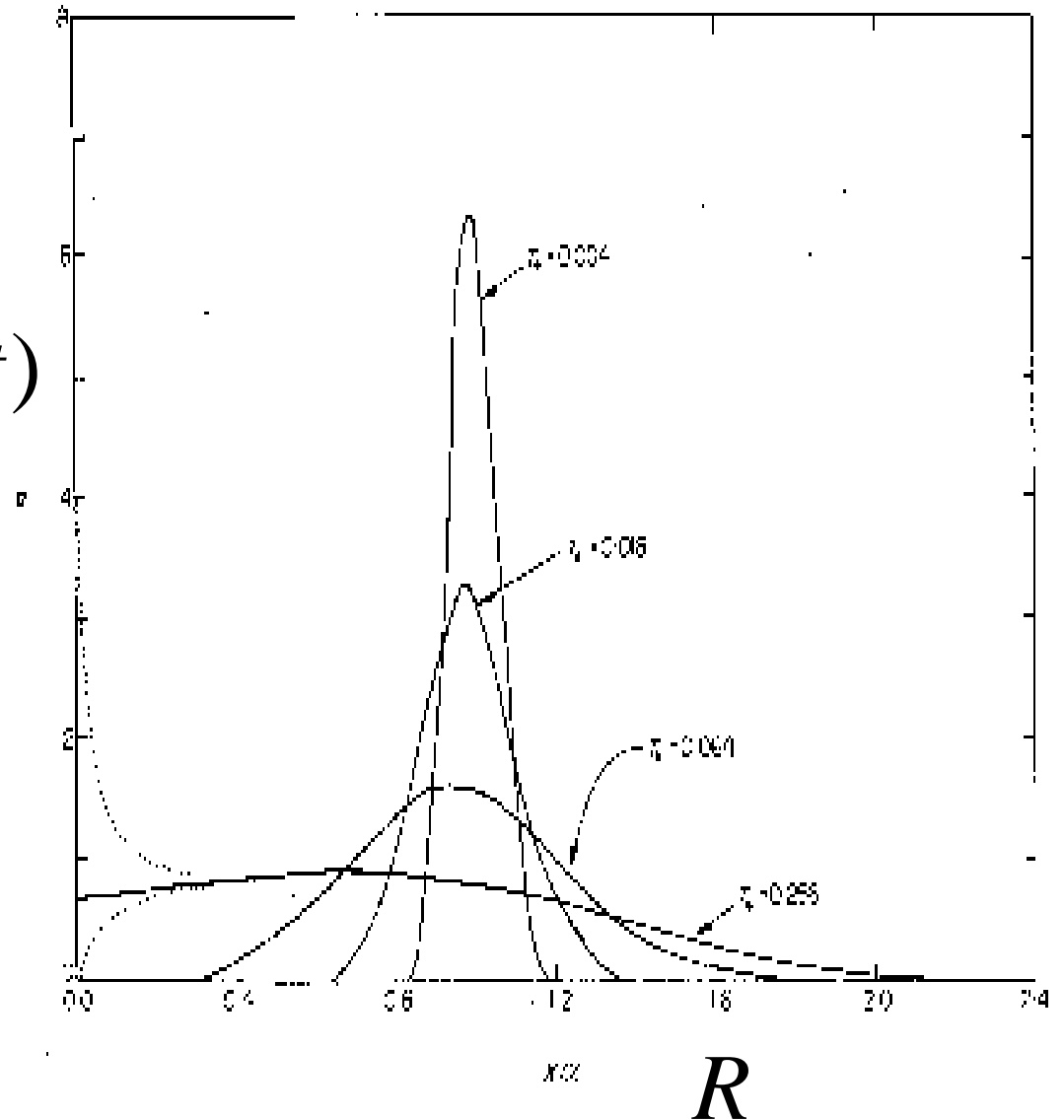
Ring spreads to form a disc

Pringle, Lynden-Bell 1974
MNRAS 168 603

$$\Sigma(R, t)$$

Surface density
evolution

diffusion in radius



Accretion Discs

- **Keplerian Orbits**

$$V_f = V_K = \left(\frac{G M}{R} \right)^{1/2} \quad \Omega \equiv \frac{V_f}{R} = \sqrt{G M R^3}$$

- **Vertical thickness: thin <--> supersonic**

$$\frac{H}{R} \approx \frac{c_s}{V_f}$$

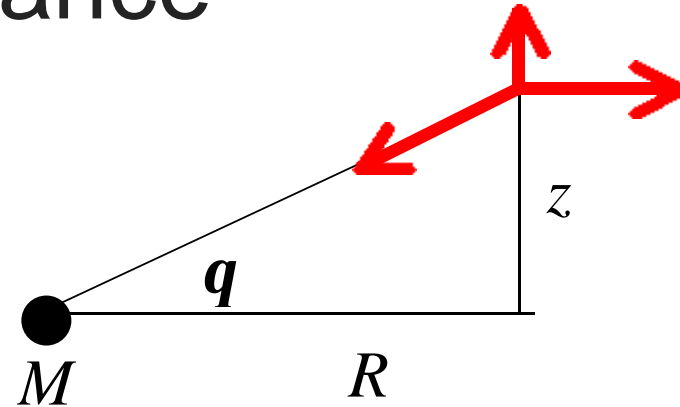
- **sound speed for perfect gas**

$$c_s^2 \equiv \frac{dP}{dr} = \mathbf{g} \frac{P}{\mathbf{r}} = \frac{\mathbf{g} k T}{\mathbf{m} m_H}$$

Force Balance

- **horizontal:**

- gravity in
- centrifugal out



$$\frac{V^2}{R} = \frac{G M}{R^2 + z^2} \cos \mathbf{q} = \frac{G M R}{(R^2 + z^2)^{3/2}} \approx \frac{G M}{R^2} \quad (z \ll R)$$

- **vertical:**

- gravity down
- pressure gradient up

$$\frac{dP}{dz} = -\mathbf{r} g_z$$

$$g_z = \frac{G M}{R^2 + z^2} \sin \mathbf{q} = \frac{G M z}{(R^2 + z^2)^{3/2}} \approx \frac{G M z}{R^3}$$

Vertical Hydrostatic Equilibrium

- Assume vertically isothermal
- vertical structure (Gaussian if iso-thermal)

$$\frac{dP}{dz} = -\mathbf{r} g_z = -\frac{\mathbf{g} P G M z}{c_s^2 R^3} = -\frac{P z}{H^2}$$

$$\frac{dP}{P} = -\frac{z dz}{H^2}$$

$$\ln P = \ln P_0 - \frac{1}{2} \left(\frac{z}{H} \right)^2$$

$$P = P_0 \exp \left\{ -\frac{1}{2} \left(\frac{z}{H} \right)^2 \right\}$$

$$H^2 = \frac{c_s^2 R^3}{\mathbf{g} G M} = \frac{k T r^3}{\mathbf{m} m_H G M}$$

$$\frac{H^2}{R^2} = \frac{1}{\mathbf{g}} \frac{c_s^2}{V_f^2}$$