



The tidal stability of Fornax cluster dwarf galaxies in Newtonian and Milgromian dynamics

Authors: Elena Asencio, Indranil Banik, Steffen Mieske, Aku Venhola, Pavel Kroupa & Hongsheng Zhao



University of St Andrews

The Fornax Deep Survey Dwarf galaxy Catalog

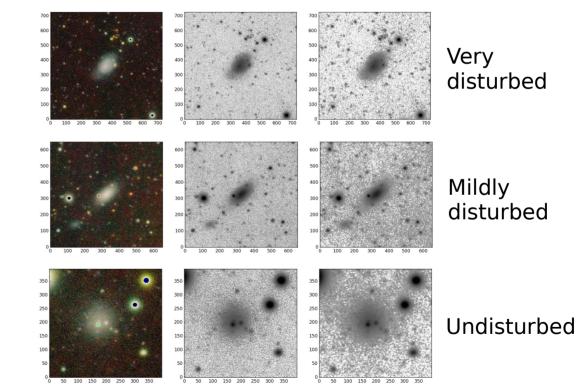
Fornax cluster:

- Second nearest galaxy cluster to us (20 Mpc away)
- Contains dwarfs with different masses and shapes

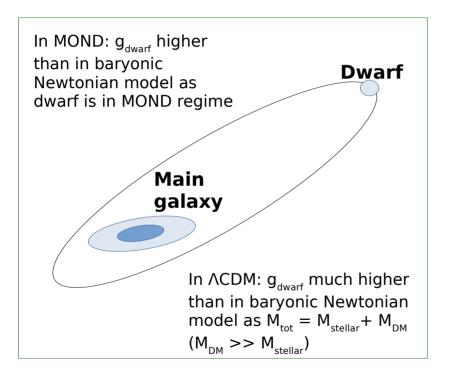
The Fornax Deep Survey Dwarf galaxy Catalog

Fornax cluster:

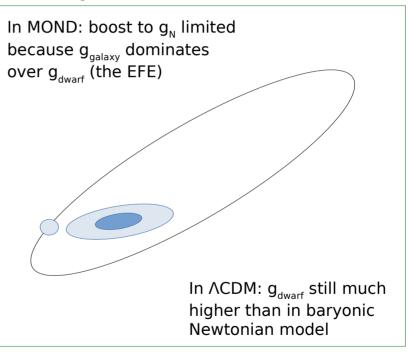
- Second nearest galaxy cluster to us (20 Mpc away)
- Contains dwarfs with different masses and shapes
- The FDSDC catalog contains 564 dwarf galaxies (353 used for the analysis)
- Most dwarf galaxies in the catalog are dE and dSph (classified as the same type)
- 50% completeness limit at $M_{_{r'}}$ = -10.5 mag (m_{_{r'}} = 21 mag) and $\mu_{_{e,r'}}$ = 26 mag arcsec^{-2}



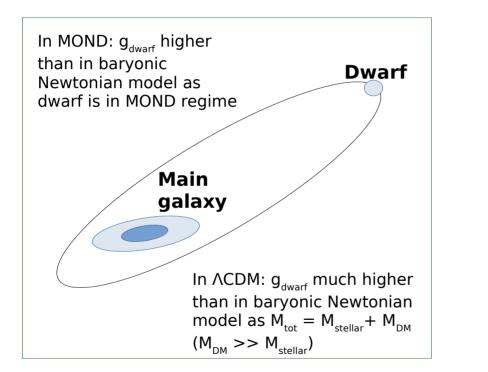
Dwarf galaxies to test gravity models



Brada & Milgrom 2000

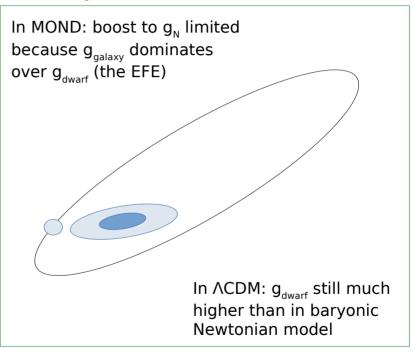


Dwarf galaxies to test gravity models



5

Brada & Milgrom 2000



Dwarf galaxies will be more disturbed by tides in MOND than in ACDM

"Note that the inner region of a satellite [in ACDM] is affected by tides after significant tidal destruction of its outer parts" (Kazantzidis et al. 2004)

Effects of gravitational interactions on dwarfs

0. Ram-pressure stripping: gas should have already been pressure stripped (Venhola+ 2019)

Effects of gravitational interactions on dwarfs

0. Ram-pressure stripping: gas should have already been pressure stripped (Venhola+ 2019)

1. Harassment: disruption due to interactions with massive galaxies

Disruption timescale: $t_d = \frac{0.043}{W} \frac{\sqrt{2} m_s r_{h,p}^2}{G m_p^2 n_p r_{h,s}^3}$ ACDM: $m = m_{stellar} + m_{DM}$ MOND: $G \rightarrow G_{eff} = G (a_0 + g_c) / g_c$ Cluster gravity

Effects of gravitational interactions on dwarfs

0. Ram-pressure stripping: gas should have already been pressure stripped (Venhola+ 2019)

Disruption timescale: $t_d = \frac{0.043}{W} \frac{\sqrt{2} m_s r_{h,p}^2}{G m_p^2 n_p r_{h,s}^3}$ ACDM: $m = m_{stellar} + m_{DM}$ MOND: $G \rightarrow G_{eff} = G (a_0 + g_c) / \underline{g_c}$

1. Harassment: disruption due to interactions with massive galaxies

2. Tidal disruption: disruption from the cluster's tidal field

$$\Lambda \text{CDM:} \quad r_{tid} = \left(\frac{G \ m_{dwarf, stellar+DM}}{2 \ (\Delta \ g_c / \Delta \ D)}\right)^{1/3} \qquad \text{Baumgardt \&} \\ \text{Makino 2003}$$

Assume 4% of total DM halo within optical
$$r_h$$
 (Díaz-García+ 2016)

MOND:
$$r_{tid} = 0.374 \left(\frac{G_{eff} m_{dwarf}}{(\Delta g_c / \Delta D)} \right)^{1/3}$$
 Zhao 2005
Zhao & Tian 2006

*We obtain
$$r_{tid}$$
 at pericentre for $P_e \propto e$: $R_{per} = 0.29 R_{3D}$ (Baumgardt priv. comm.)

Tidal stress (observed

from X-rays)

 $\frac{G \ m_{dwarf}}{r_{tid}^2} \approx r_{tid} \frac{\Delta \ g_c}{\Delta \ D}$

Cluster

gravity

Tidal susceptibility (ŋ)

• Tidal susceptibility from harassment:

 $\eta_{\scriptscriptstyle har} \equiv t_{\scriptscriptstyle Fornax} / t_d$

With $t_{Formax} = 10 \pm 1 \text{ Gyr}$ (Rakos+ 2001)

If $t_d >> t_{Fornax}$: η_{har} very small (the dwarf will **not** be very affected by harassment)

If $t_d << t_{Fornax}$: η_{har} very high (the dwarf will be very affected by harassment)

Tidal susceptibility (ŋ)

• Tidal susceptibility from harassment:

 $\eta_{har} \equiv t_{Fornax} / t_d$

With $t_{Fornax} = 10 \pm 1 \text{ Gyr}$ (Rakos+ 2001)

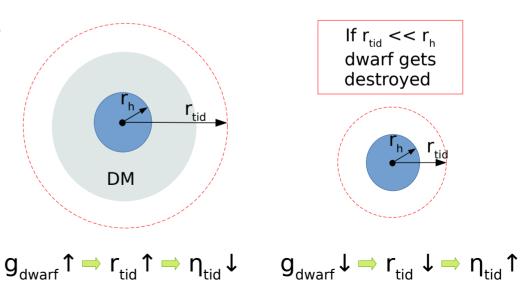
• Tidal susceptibility from cluster tidal field at pericentre:

 $\eta_{\scriptscriptstyle tid} \equiv r_{\scriptscriptstyle h} / r_{\scriptscriptstyle tid}$

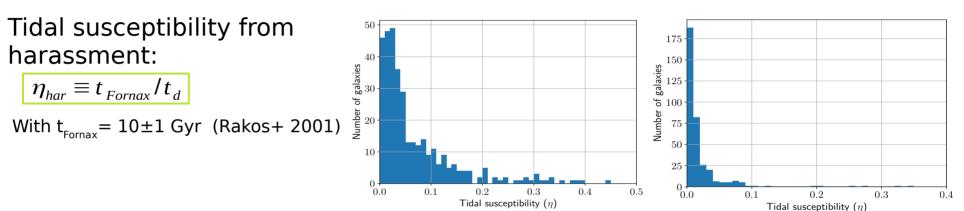
 $r_h \equiv$ radius containing half of the total luminous mass of the object

 $r_{tid} \equiv$ radius at which the gravitational tide from an external object starts to dominate over the self-gravity of the object If $t_d >> t_{Formax}$: η_{har} very small (the dwarf will **not** be very affected by harassment)

If $t_d << t_{Fornax}$: η_{har} very high (the dwarf will be very affected by harassment)



Tidal susceptibility (η) values



ΛCDM

MOND

•

Tidal susceptibility (η) values

• Tidal susceptibility from harassment:

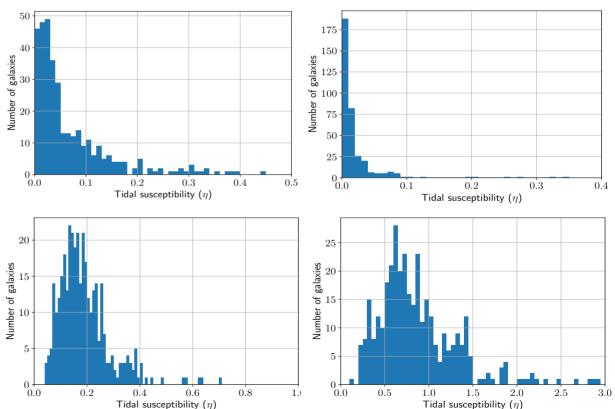
 $\eta_{har} \equiv t_{Fornax} / t_d$

With $t_{Fornax} = 10 \pm 1 \text{ Gyr}$ (Rakos+ 2001)

• Tidal susceptibility from cluster tidal field:

 $\eta_{tid} \equiv r_h / r_{tid}$

With $r_h \approx \frac{4}{3} R_e$ (Baumgardt+ 2010)



ΛCDM

Tidal susceptibility (n) values

Tidal susceptibility from harassment:

 $\eta_{har} \equiv t_{Fornax}/t_d$

With $t_{Formax} = 10 \pm 1$ Gyr (Rakos+ 2001)

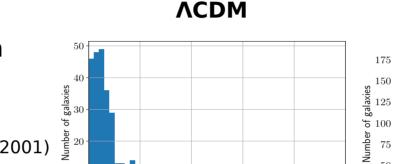
Tidal susceptibility from ٠ cluster tidal field:

 $\eta_{tid} \equiv r_h / r_{tid}$

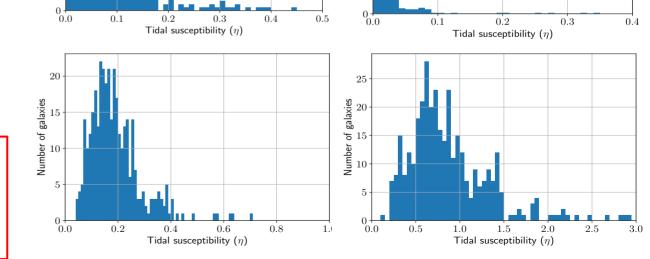
Fk

With $r_h \approx \frac{4}{3} R_e$ (Baumgardt+ 2010)

- Effect of η_{har} is negligible in both cosmologies
- MOND η_{tid} is about 5x higher • than in ΛCDM .



 10^{-1}



175

150

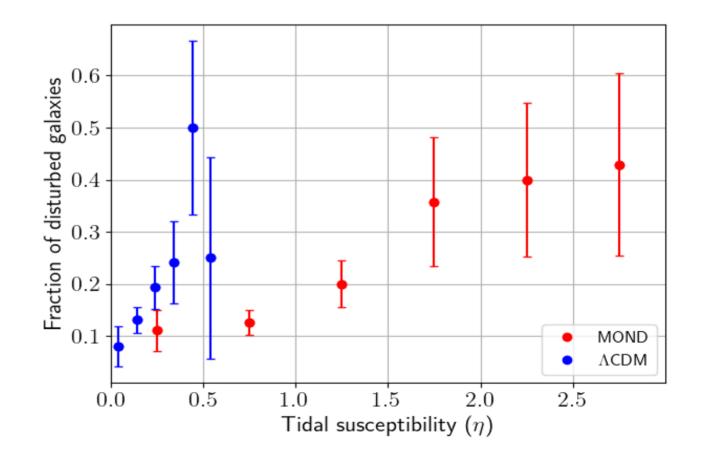
7550

25



Testing the models with disturbed fraction vs η

- > We expect that Fornax dwarfs with $\eta > (0.5 - 1)$ will be tidally disturbed
 - ACDM: trend goes up at η significantly lower than expected
 - Lack of dwarfs that should still be tidally stable
 - MOND: trend goes up at η a bit higher than expected



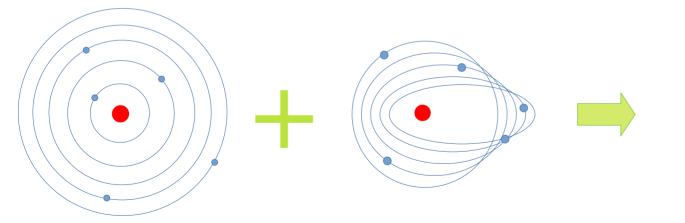
Comparing n with morphological classification

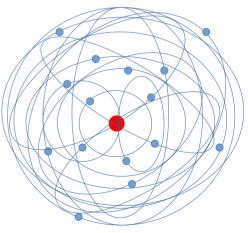
- At which η do the dwarfs start being classified as "perturbed" (in the catalogue) in each model? \implies find min η_{dist} value
- What is the maximum η reached by the dwarfs before being destroyed in each model ? \implies find η_{destr} value

Test particle simulation: step 1 (grid of orbits)

Step 1: We simulate orbits of test masses in the observed cluster potential for a grid with all possible distance (R_i) and eccentricity (e) values.

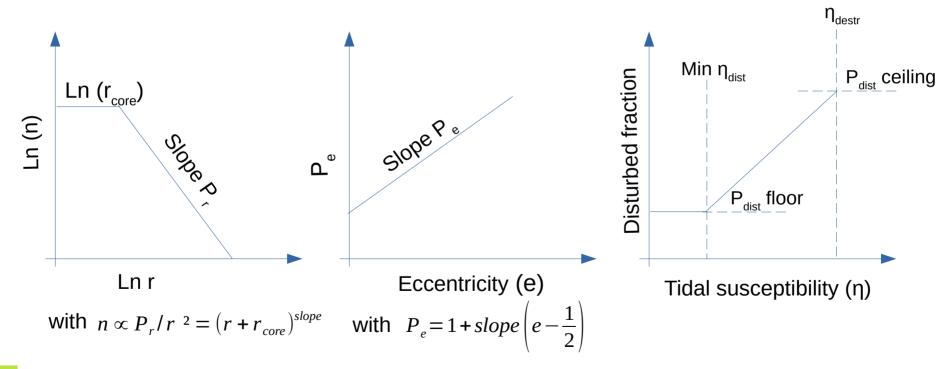
- Record max η over the orbit, use it to assign disturbed probability or destruction (next slide)
- Consider sky-projected separation from all possible angles.





Test particle simulation: step 2 (statistics)

Step 2: assign probabilities to the orbits:



Test particle simulation: step 2 (statistics)

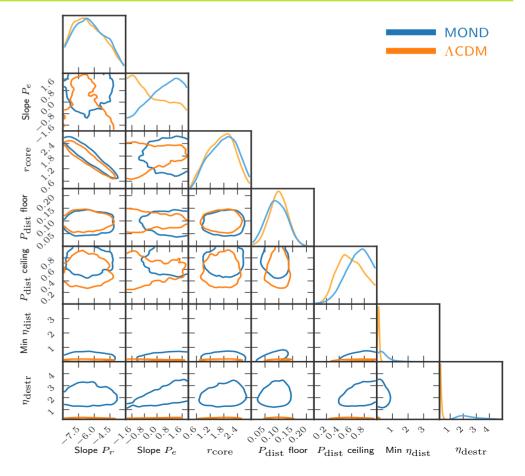
Step 2: assign probabilities to the orbits:

In order to fit the observational data we leave as free parameters:

- r_{core} : radius of the constant density central region of the Fornax cluster
- Slope P_r : power-law slope of dwarf radial distribution in cluster outskirts $P_r = r^2 (r + r_{core})^{slope}$
- Slope P_e : slope of the eccentricity probability distribution $P_e = 1 + slope \left(e \frac{1}{2}\right)$
- Min η_{dist} : lowest η value at which the dwarf is disturbed.
- η_{destr} : η value at which the dwarf is destroyed.
- P_{dist} floor: minimum probability for a dwarf to appear disturbed if $\eta < \min \eta_{dist}$ (e.g. due to asymmetric star formation)
- P_{dist} ceiling: probability for a dwarf to appear disturbed right before it gets destroyed ($\eta = \eta_{destr}$)

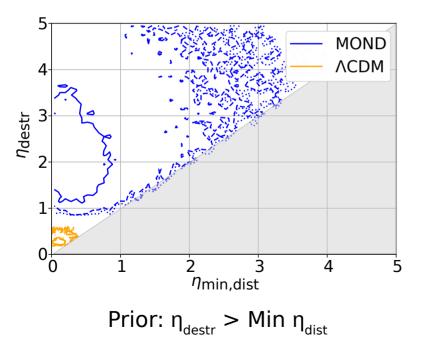
Step 4: finding the best model and uncertainties

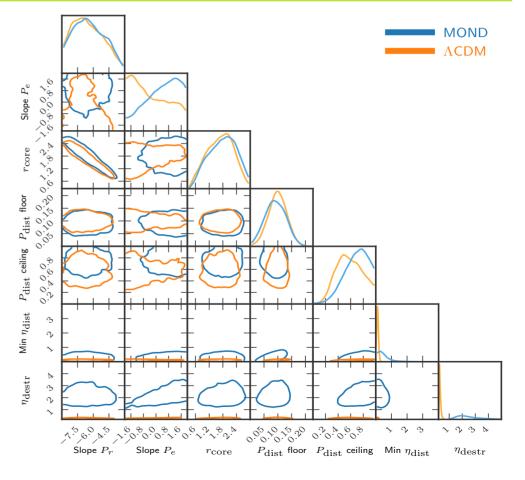
Step 4: to find the set of simulation parameter values that provide a good match to the observed population, we use the Markov chain Monte Carlo (MCMC) method.



Step 4: finding the best model and uncertainties

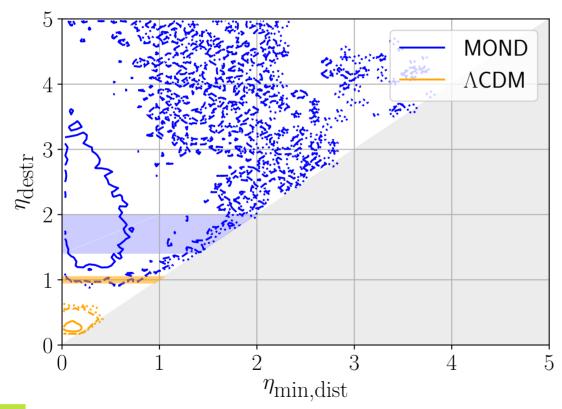
Step 4: to find the set of simulation parameter values that provide a good match to the observed population, we use the Markov chain Monte Carlo (MCMC) method.





At which n are the dwarfs actually expected to be destroyed/severely perturbed?

Interpreting MCMC results with *N*-body models



ACDM N-body:

$$\eta_{\rm destr} pprox 1$$

(Peñarrubia+ 2009)

MOND N-body:

 $\eta_{destr} = 1.7 \pm 0.3$

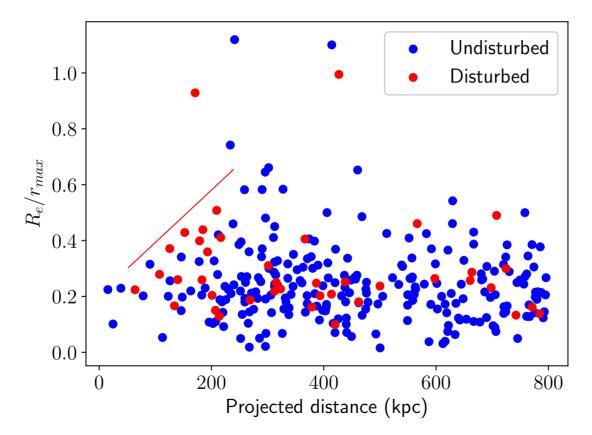
(Asencio+ 2022)

Conclusions

- Observations of Fornax dwarf morphologies tell us that some are disturbed
 - Disturbed fraction higher towards center
- Main process expected to be cluster tides (dwarfs should be gas poor)
 - + We expect max η (r_{_h} / r_ $_{tidal})~\approx 1$
- ACDM: Fornax dwarfs should **not** be tidally disturbed
 - But observations imply they are disturbed (not due to detection limit)
 - This requires stability limit of $\eta_{destr} = 0.25^{+0.07}_{-0.03}$ to match observations (by 10⁵ MCMC trials)
 - (Tidal force)/(Internal gravity) $\approx \eta^3$
- MOND: Fornax dwarfs ${\bm are}$ expected to be disturbed (η is higher in this model due to EFE and lack of cold dark matter)
 - The required stability limit is $\eta_{destr} = 1.88^{+0.85}_{-0.53}$
 - *N*-body simulations imply $\eta_{destr} = 1.7 \pm 0.3$

Appendix

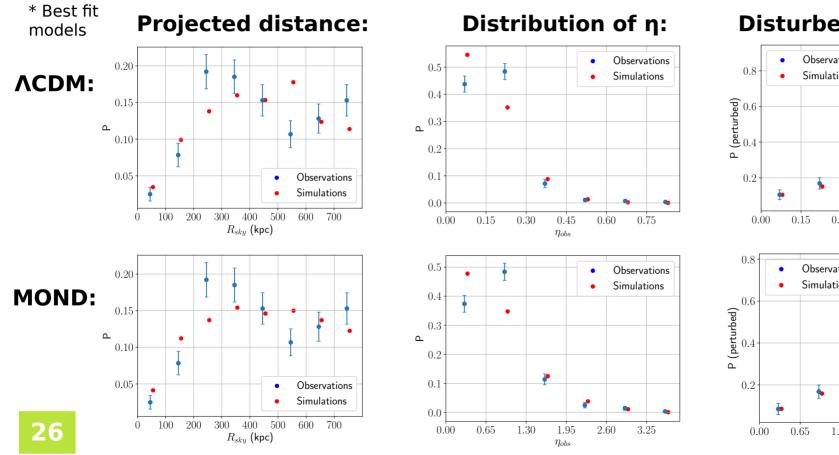
Tides remove dwarfs close to cluster center



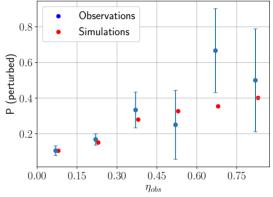
- Dwarfs with larger size at fixed mass are more susceptible to tides, but also harder to detect
- However, selection effects alone insufficient to explain lack of diffuse dwarfs towards cluster center (above red line)
- Most disturbed dwarfs at projected distance < 500 kpc from the center.

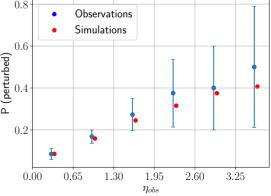
 $* r_{max}$ is the maximum R_e at fixed stellar mass for the dwarf to remain detectable

Step 3: observational constraints

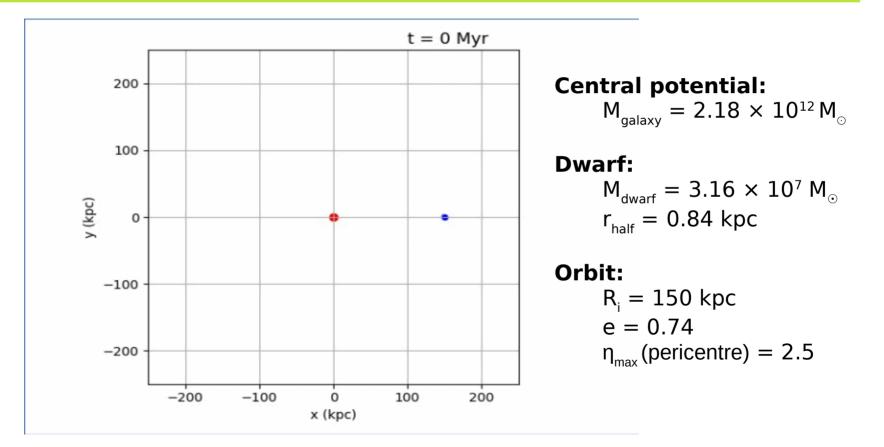


Disturbed fraction vs η:



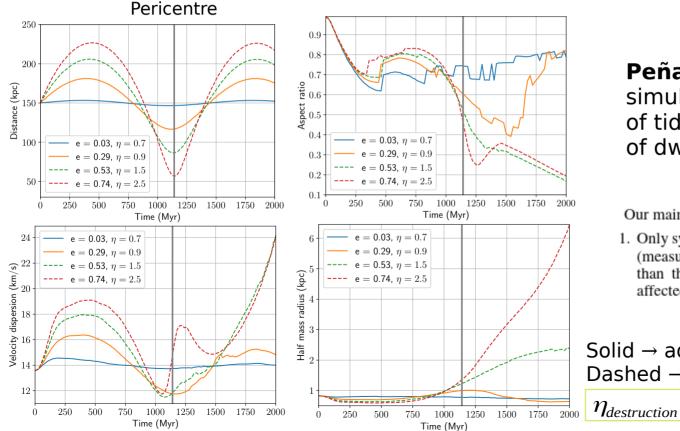


N-body simulations in MOND (Nagesh+ 2021)



N-body simulations

MOND



ΛCDM

Peñarrubia+ 2009: *N*-body simulations to explore the effects of tidal stripping on the structure of dwarf spheroidal galaxies

Our main findings may be summarized as follows.

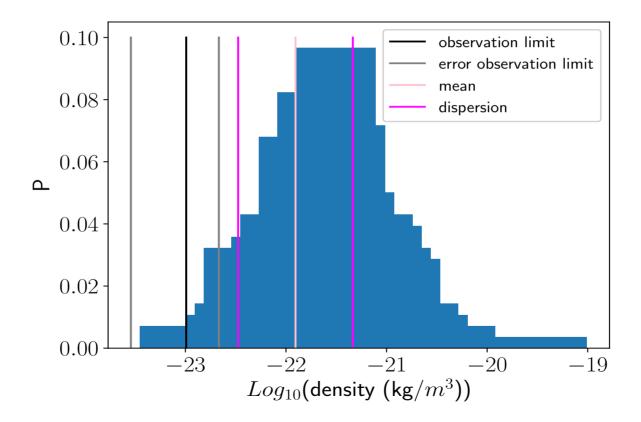
 Only systems in orbits where the tidal radius (Equation (3)) (measured at perigalacticon) is comparable to or smaller than the luminous radius of the dwarf are significantly affected by tides.

Solid \rightarrow adiabatic response Dashed \rightarrow destroyed $n_{\rm destroyed} \approx 1.5$

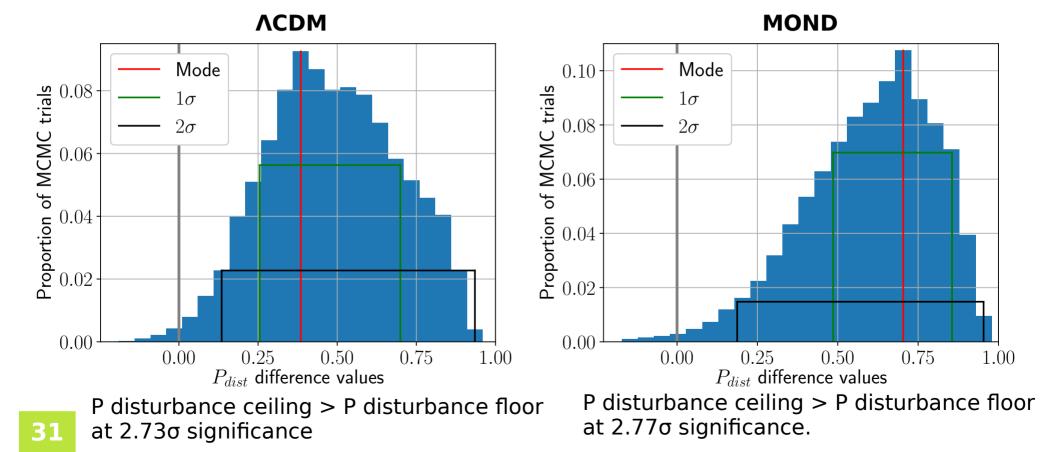
Cosmological MOND framework (vHDM)

- Proposed by Angus 2009 (MNRAS, 394, 527)
- Cold dark matter (CDM) replaced by fast collisionless matter
 - e.g. 11 eV/c² sterile neutrinos (Angus+2007)
 - Only in galaxy clusters (galaxies unaffected by neutrinos if $m_v < 100 \text{ eV/c}^2$)
- MOND is applied only to density perturbations
- MOND effects become important only at z < 50
 - e.g. Nusser 2002, Llinares+ 2008, Angus+ 2013, Katz+ 2013, Candlish 2016
- Standard background cosmology, expansion and thermal history
- It can explain:
 - BBN
 - CMB
 - Bullet Cluster and 30 virialized clusters (Angus+ 2010, MNRAS, 402, 395)
 - Problems with ACDM on galaxy scales (e.g: planes of satellites problem)
 - KBC void and Hubble tension (MNRAS, 499, 2845)
 - El Gordo galaxy cluster (MNRAS, 500, 5249)

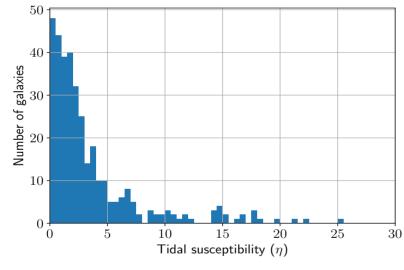
Density distribution of the dwarfs within r_h



Histograms (P disturbance ceiling - P disturbance floor)

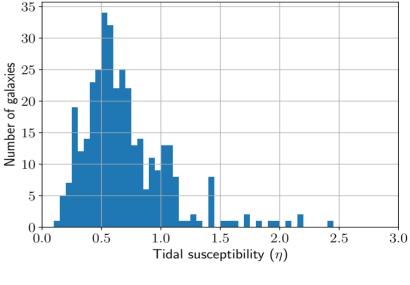


Tidal susceptibility in baryonic Newtonian model



 $\eta_{har} \equiv t_{Fornax} / t_{d}$

Harassment



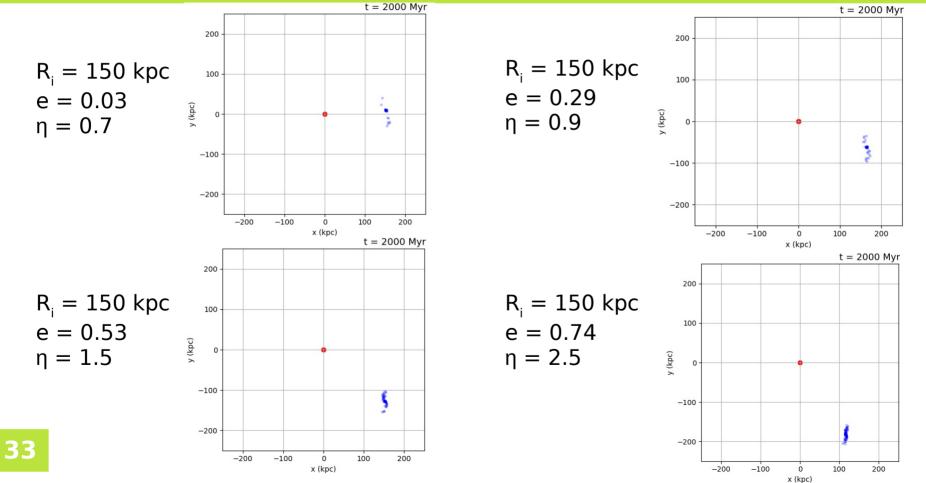
 $\eta_{tid} \equiv r_h / r_{tid}$

Cluster tides

of ga

32

Final snapshots for dwarfs with different η



 $R_i \equiv initial R$ = semimajor axis

LOCAL GROUP DWARF SPHEROIDALS: CORRELATED DEVIATIONS FROM THE BARYONIC TULLY-FISHER RELATION

THE ASTROPHYSICAL JOURNAL, 722:248-261, 2010 October 10 © 2010. The American Astronomical Society. All rights reserved. Printed in the U.S.A. doi:10.1088/0004-637X/722/1/248

LOCAL GROUP DWARF SPHEROIDALS: CORRELATED DEVIATIONS FROM THE BARYONIC TULLY–FISHER RELATION

STACY S. MCGAUGH¹ AND JOE WOLF²

¹ Department of Astronomy, University of Maryland, College Park, MD 20742-2421, USA; ssm@astro.umd.edu
² Center for Cosmology, Department of Physics and Astronomy, University of California, Irvine, CA 92697-4575, USA; wolfj@uci.edu *Received 2010 March 17; accepted 2010 August 11; published 2010 September 17*

ABSTRACT

Local Group dwarf spheroidal satellite galaxies are the faintest extragalactic stellar systems known. We examine recent data for these objects in the plane of the Baryonic Tully–Fisher Relation (BTFR). While some dwarf spheroidals adhere to the BTFR, others deviate substantially. We examine the residuals from the BTFR and find that they are not random. The residuals correlate with luminosity, size, metallicity, ellipticity, and susceptibility of the dwarfs to tidal disruption in the sense that fainter, more elliptical, and tidally more susceptible dwarfs deviate farther from the BTFR. These correlations disfavor stochastic processes and suggest a role for tidal effects. We identify a test to distinguish between ACDM and MOND based on the orbits of the dwarf satellites of the Milky Way and how stars are lost from them.

Key words: dark matter - galaxies: dwarf - galaxies: formation - galaxies: halos - Local Group

Online-only material: color figures

34

LOCAL GROUP DWARF SPHEROIDALS: CORRELATED DEVIATIONS FROM THE BARYONIC TULLY-FISHER RELATION

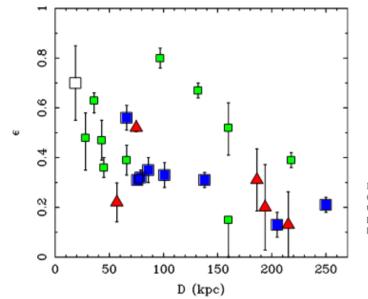


Figure 5. Ellipticity of Local Group dwarf spheroidals as a function of distance from their host galaxy. Symbols as per Figure 1. The open square is Sagittarius (Ibata et al. 1997). Leo T ($\epsilon = 0.29$ at D = 407 kpc) falls off the right edge of the plot.

- Non-circular photos of Milky Way satellites suggest tidal disturbance
- Ellipticities higher at lower distance, as expected for tides

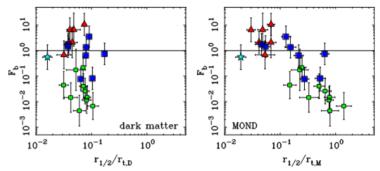


Figure 6. Residuals from the BTFR (symbols as per Figure 1) plotted against the ratio of half-mass radius to tidal radius in the case of dark matter (left) and MOND (right). The plots are scaled identically. The tidal radius depends on the masses of satellite and host (Equations (4) and (5)). Masses are computed dynamically for the case of dark matter (left: $m = M_{1/2}$) and using only baryonic mass in the case of MOND (right: $m = M_b$). The arrow marks the location where the average photometric tidal radius equals the computed tidal radius. In the case of MOND, the location where dwarfs deviate form the BTFR corresponds well to the observed photometric tidal radius, and the amount of deviation correlates with the size of a dwarf relative to its MONDian tidal radius.

- F_b is measure of discrepancy from isolated MOND prediction ($F_b < 1$: $\sigma_{obs} > \sigma_{MOND}$)
- MOND predictions assuming virial equilibrium do not work in many cases, but tides expected to be significant in these cases
- MOND works well when satellite expected to be tidally stable (low η)
- ACDM predicts no tidal disturbance in all cases, which may conflict with observed signs of disturbance.