

Importance of Accretion Discs

- **Formation of compact objects**
 - friction moves angular momentum outward
 - allows matter to spiral inward
 - build up compact object at centre
- **Generation of light**
 - gravitational potential energy
 - converted by friction to heat
 - radiated as light

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3 Classic Papers

- **Black hole accretion discs**
Shakura, Sunyaev 1973 A&A 24 337
- **Time-dependent discs**
Pringle, Lynden-Bell 1974 MNRAS 168 603
- **Gas streams**
Lubow, Shu 1975 ApJ 198 383

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Accretion Energy

- **Nuclear energy (H -> He) Efficiency:**

$$\frac{\Delta E_{nuc}}{m} \approx 6 \times 10^{14} \frac{\text{J}}{\text{kg}} \quad \mathbf{h} = \frac{\Delta E_{nuc}}{m c^2} = 0.7\%$$
- **Accretion energy**

$$\frac{\Delta E_{acc}}{m} \approx \frac{GM}{R} \approx 1.7 \times 10^{11} \left(\frac{M}{M_{sun}} \right) \left(\frac{R}{R_{sun}} \right)^{-1} \frac{\text{J}}{\text{kg}}$$

$$\approx 1.7 \times 10^{16} \left(\frac{M}{1.4 M_{sun}} \right) \left(\frac{R}{10 \text{ km}} \right)^{-1} \frac{\text{J}}{\text{kg}}$$
- **more efficient for compact objects**

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Efficiency

- **Nuclear energy (H -> He) Efficiency:**

$$\Delta E_{nuc} = \mathbf{h} m c^2 \quad \mathbf{h} = 0.7\%$$
- **Accretion efficiency**

$$L \approx \frac{GM \dot{M}}{R} = \mathbf{h} \dot{M} c^2 \quad \mathbf{h} = \frac{GM}{R c^2} = \frac{R_H}{2R} \quad R_H \approx 3 \text{ km} \left(\frac{M}{M_{sun}} \right)$$
- **Compactness --- M / R**
- **Black hole - smallest stable orbit**
 - non-rotating $R = 6 R_H \quad \mathbf{h} = \frac{1}{12} \approx 8\%$
 - maximally-rotating $R = 3 R_H \quad \mathbf{h} = \frac{1}{2} \approx 17\%$
- **White Dwarf**

$$M \approx M_{sun} \quad R \approx 10^4 \text{ km} \quad \mathbf{h} \sim 10^{-4}$$

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Radiation Opposes Gravity

- **Coulomb attraction ties electrons and protons**
- **gravity pulls in**

$$F_{grav} = -\frac{GMm}{r^2}$$
- **radiation pushes electrons out**

$$F_{rad} = \frac{L \sigma_T}{4\pi r^2} \quad L_{acc} = \frac{GM\dot{M}}{R}$$

Thompson electron scattering cross-section $\sigma_T \approx 6.65 \times 10^{-29} \text{ m}^2$
- **Notes:**
 - radiation pressure opposes gravity
 - same R scaling
 - radiation pressure > gravity at high \dot{M}

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Eddington Luminosity

- **Photon momentum** $p_g = \frac{h\mathbf{n}}{c}$
- **Density of photons** $n_g = \frac{L/h\mathbf{n}}{4\mathbf{p} R^2 c}$
- **Radiative force (per electron)** $F_{rad} = n_g p_g s_T c = \frac{L s_T}{4\mathbf{p} R^2 c}$ $s_T = 6.7 \times 10^{-25} \text{ cm}^2$
- **Total force (per electron + proton pair)** $F_{rad} + F_{grav} = \frac{L s_T}{4\mathbf{p} R^2 c} - \frac{GM(m_p + m_e)}{R^2}$
- **Eddington Luminosity** $L_{Edd} = \frac{4\mathbf{p} G M (m_p + m_e) c}{s_T} = 1.3 \times 10^{38} \left(\frac{M}{M_{sun}} \right) \text{ erg s}^{-1}$

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Eddington Accretion Rate

- **Eddington Luminosity** $L_{Edd} = \frac{4\mathbf{p} G M m_p c}{s_T}$
- **Accretion Luminosity** $L_{acc} = \frac{G M \dot{M}}{R}$
- **Eddington Accretion Rate** $\frac{L_{acc}}{L_{Edd}} = \frac{s_T \dot{M}}{4\mathbf{p} m_p c R} = \frac{\dot{M}}{\dot{M}_{Edd}}$

$$\dot{M}_{Edd} = \frac{4\mathbf{p} m_p c R}{s_T} \approx 10^{-5} \left(\frac{R}{R_{sun}} \right) \frac{M_{sun}}{\text{yr}}$$

	kg/s	M_{sun}/yr
WD	6×10^{19}	10^{-7}
NS	10^{12}	1.4×10^{-12}

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Rough Temperatures

- **Optically thick**
 - Blackbody radiates accretion luminosity
$$s T_b^4 = \frac{L_{acc}}{4\mathbf{p} R^2} \quad T_b = \left(\frac{L_{acc}}{4\mathbf{p} R^2 s} \right)^{1/4} \approx \left(\frac{G M \dot{M}}{4\mathbf{p} R^2 s} \right)^{1/4}$$
- **Optically thin**
 - potential energy released
 - = thermal energy of shocked gas
$$\frac{G M m_p}{R} = 2 \times \frac{3}{2} k T_{th} \quad T_{th} \approx \frac{G M m_p}{3 k R}$$
- **Radiation temperature and photon energy**

$$T_b < T_{rad} < T_{th} \quad k T_b < h\mathbf{n} < k T_{th}$$

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Rough Temperatures

- **Neutron Star or Black Hole**

$$L_{acc} \sim L_{Edd} \sim 10^{38} \text{ ergs}^{-1}$$

$$10^7 \text{ K} < T_{rad} < 10^{11} \text{ K}$$

$$1 \text{ keV} < h\mathbf{n} < 50 \text{ MeV}$$
 - mid to hard X-rays
- **White Dwarf**

$$L_{acc} \sim 10^{33} \text{ ergs}^{-1}$$

$$6 \times 10^4 \text{ K} < T_{rad} < 10^9 \text{ K}$$

$$6 \text{ eV} < h\mathbf{n} < 100 \text{ keV}$$
 - optical - uv - soft X-ray

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Roche Lobe Overflow

- **Initial stream velocity** $V_{\perp} \sim c_s \sim 10 \text{ km s}^{-1}$
- **L1 velocity relative to centre of mass**

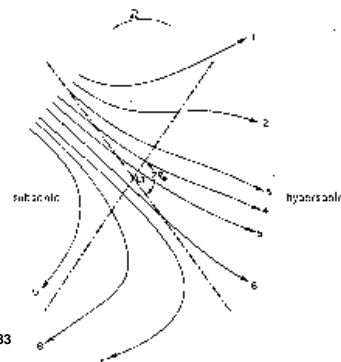
$$V_{\perp} = \frac{2\mathbf{p}}{P} R_{L1} \approx \frac{2\mathbf{p} a}{P} (0.5 - 0.23 \log q)$$

$$\sim 100 \text{ km s}^{-1} \left(\frac{M}{M_{sun}} \right)^{1/3} \left(\frac{P}{\text{day}} \right)^{-1/3}$$
- **subsonic --> supersonic transition at nozzle**
- **ballistic trajectory in Roche potential**
 - (neglect pressure forces)

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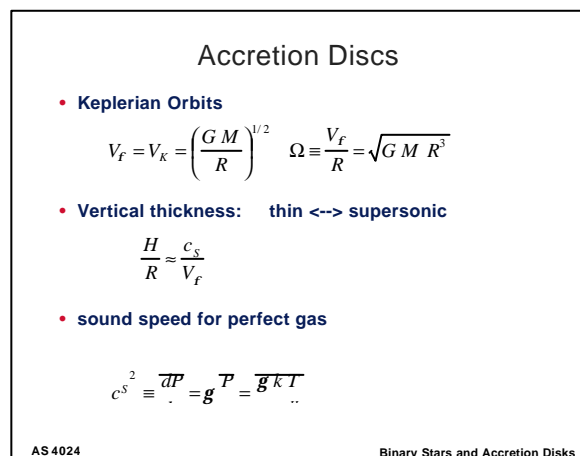
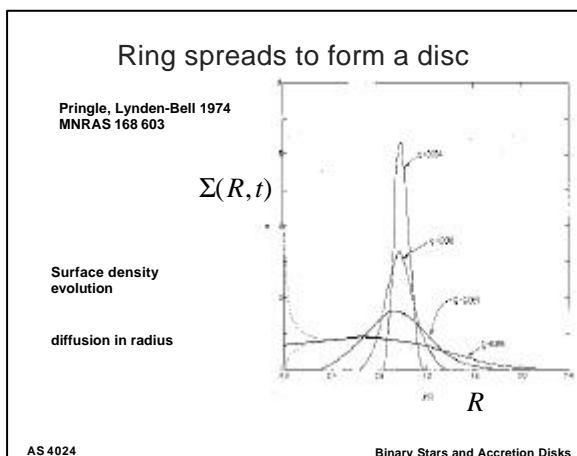
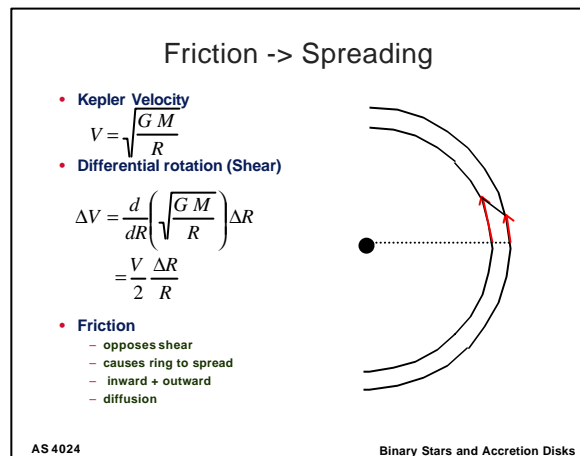
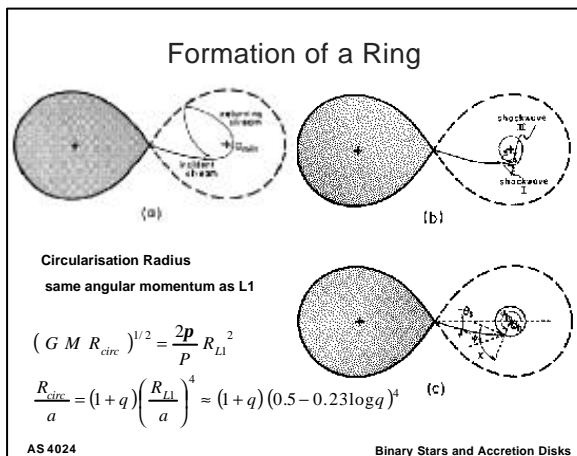
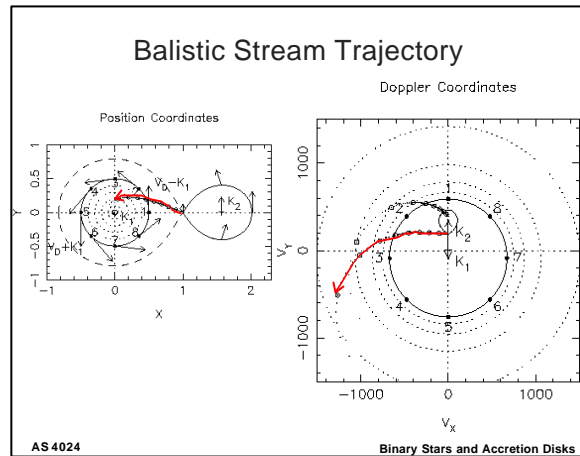
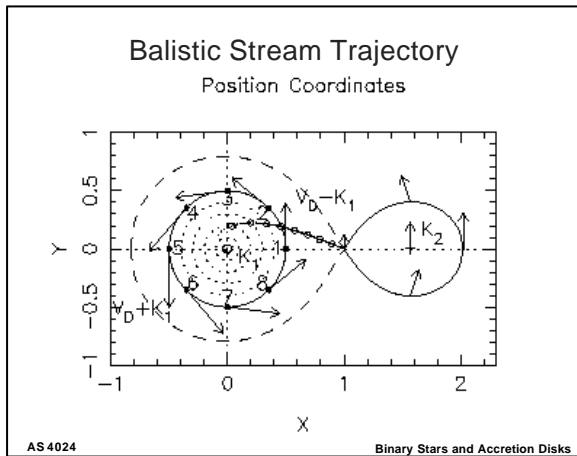
Flow thru the L1 nozzle

subsonic -> hypersonic



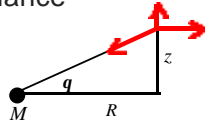
Lubow, Shu 1975 ApJ 198 383

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Force Balance

- **horizontal:**
 - gravity in
 - centrifugal out



$$\frac{V^2}{R} = \frac{GM}{R^2 + z^2} \cos q = \frac{GM R}{(R^2 + z^2)^{3/2}} \approx \frac{GM}{R^2} \quad (z \ll R)$$

- **vertical:**
 - gravity down
 - pressure gradient up

$$\frac{dP}{dz} = -r g_z$$

$$g_z = \frac{GM}{R^2 + z^2} \sin q = \frac{GM z}{(R^2 + z^2)^{3/2}} \approx \frac{GM z}{R^3}$$

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Vertical Hydrostatic Equilibrium

- Assume vertically isothermal
- vertical structure (Gaussian if iso-thermal)

$$\frac{dP}{dz} = -r g_z = -\frac{g P}{c_s^2} \frac{GM z}{R^3} = -\frac{P z}{H^2}$$

$$\frac{dP}{P} = -\frac{z dz}{H^2}$$

$$\ln P = \ln P_0 - \frac{1}{2} \left(\frac{z}{H} \right)^2$$

$$P = P_0 \exp \left\{ -\frac{1}{2} \left(\frac{z}{H} \right)^2 \right\}$$

$$H^2 = \frac{c_s^2 R^3}{g GM} = \frac{k T r^3}{m m_H GM}$$

$$\frac{H^2}{R^2} = \frac{1}{g} \frac{c_s^2}{V_f^2}$$

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